

Přednáška 5.(29.10.2007)

Měření tvaru jader pomocí mionových atomů.

Štěpení těžkých jader (fission)

Jaderné reakce.

Fúze lehkých jader (fusion)

Objev pozitronu v kosmickém záření

Atom vodíku - opakování

$$l=0$$

$$-\frac{(\hbar c)^2}{2m_e} \frac{d^2}{dr^2} u(r) - \alpha \frac{\hbar c}{r} u(r) = E u(r)$$

$$r_B = \frac{(\hbar c)^2}{2m_e} \frac{2}{\alpha \hbar c} = \frac{\hbar c}{m_e \alpha} = \frac{197 \text{ MeV fm}}{0,511 \text{ MeV} / 137} = 52816 \text{ fm}$$

$$E = -\frac{(\hbar c)^2}{2m_e} \frac{1}{r_B^2} = -\frac{(\hbar c)^2}{2m_e} \frac{(m_e \alpha)^2}{(\hbar c)^2} = -\frac{1}{2} m_e \alpha^2 = -\frac{1}{2} 0,511 \text{ MeV} \frac{1}{137^2} = -13,6 \text{ eV}$$

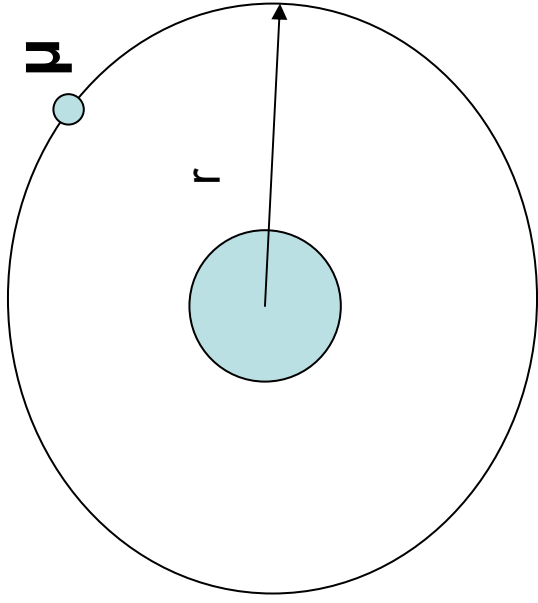
$$E = -\frac{1}{2} m_e \alpha^2 ; E = T + V = \frac{1}{2} m_e \alpha^2 - 2 \frac{1}{2} m_e \alpha^2 = 13,6 \text{ eV} - 27,2 \text{ eV}$$

$$T = \frac{1}{2} m_e \left(\frac{v}{c} \right)^2 = \frac{1}{2} m_e \beta^2 ; \text{ rychlost elektronu } \beta = \frac{1}{137} = \alpha$$

atom s Z protony :

$$E = -\frac{1}{2} m_e (Z\alpha)^2 = Z^2 13,6 \text{ eV} ; E_{Fe} = -26^2 13,6 \text{ eV} = -9,19 \text{ keV}$$

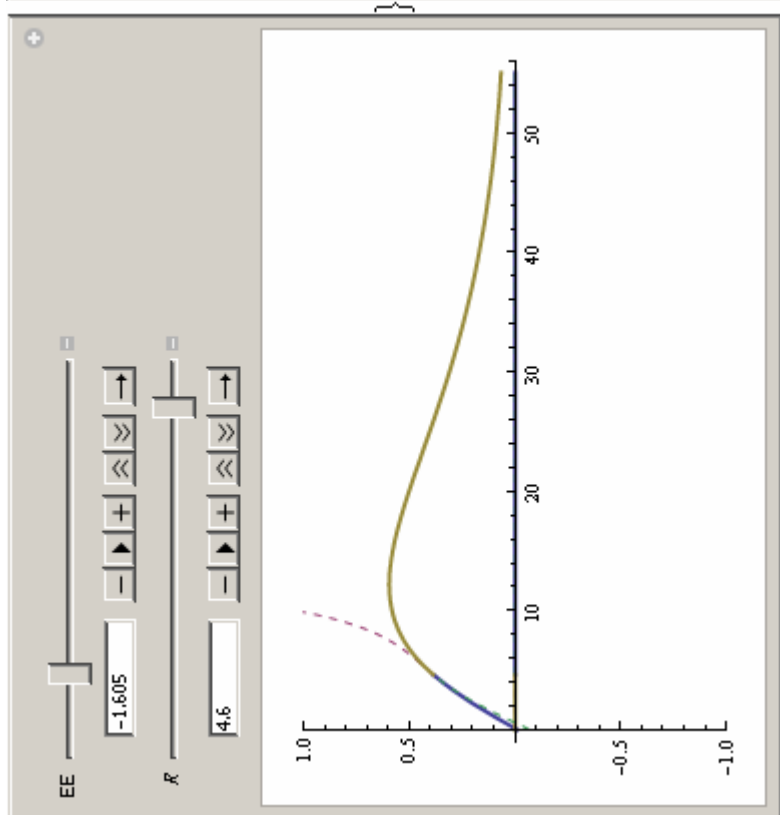
Mionový atom: elektron je nahrazen záporně nabitým mionem



$$r = r_B \frac{m_e}{m_\mu} \frac{1}{Z} \Rightarrow r = 52 \cdot 10^3 \text{ fm} \frac{0,511 \text{ MeV}}{105,6 \text{ MeV}} \frac{1}{26} \cong 9,7 \text{ fm}$$

$$R = 0 \Rightarrow E = -13,6 \text{ eV} \cdot \frac{m_\mu}{m_e} \cdot Z^2 \cdot \frac{105,6 \text{ MeV}}{0,511 \text{ MeV}} \cdot 26^2 = -1,9 \text{ MeV}$$

$$R = 1,2 \text{ fm} \cdot \sqrt[3]{A_{Fe}} = 1,2 \text{ fm} \cdot \sqrt[3]{56} \cong 4,6 \text{ fm} \Rightarrow E \cong -1,605 \text{ MeV}$$





The Nobel Prize in Chemistry 1944

"for his discovery of the fission of heavy nuclei"

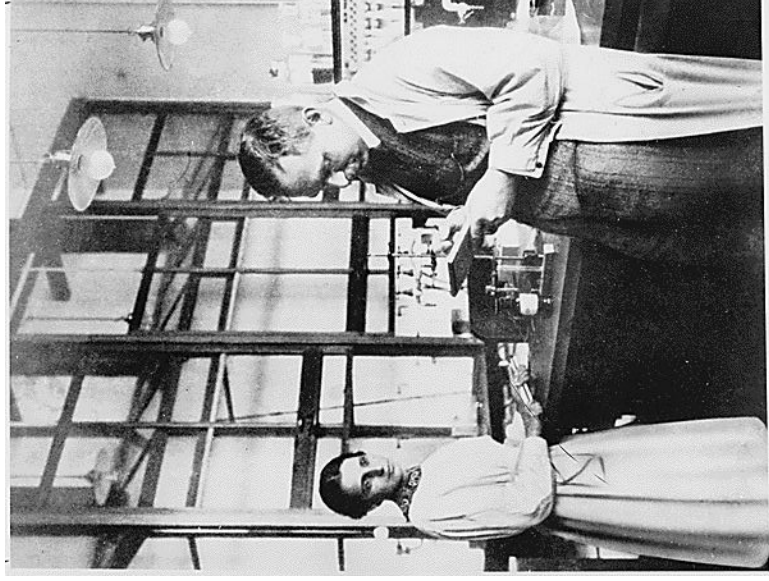


Otto Hahn

Germany

Kaiser-Wilhelm-Institut
(now Max-Planck Institut)
für Chemie
Berlin-Dahlem, Germany

b. 1879
d. 1968



O. Hahn a L. Meitner ozařovali Uran neutrony ve snaze vyrobit prvky těžší než Uran. K velkému překvapení však objevili, že se rodí Barium, tj. že se Uran štěpí na lehčí prvky.

Energie uvolněná při štěpení

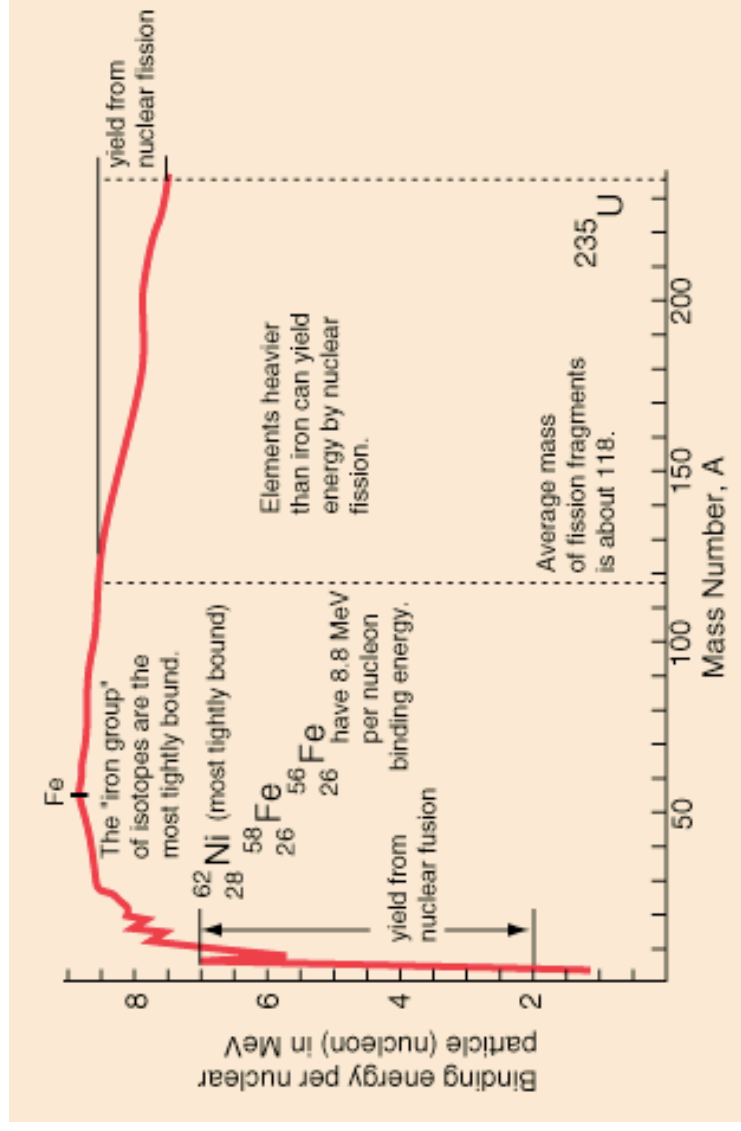
Příklad vznik Ba



$$Q = m_n + 92m_p + 143m_n - B(235,92) -$$

$$(56m_p + 84m_n - B(140,56) + 36m_p + 57m_n - B(93,36) + 3m_n) =$$

$$B(140,56) + B(93,36) - B(235,92) \simeq 160\text{MeV}$$



K tomu ještě 5MeV neutrony a 35 MeV radioaktivní rozpady celkem tedy asi 200 MeV

Vzniklá jádra mají příliš mnoho neutronů = nacházejí se mimo údolí stability.

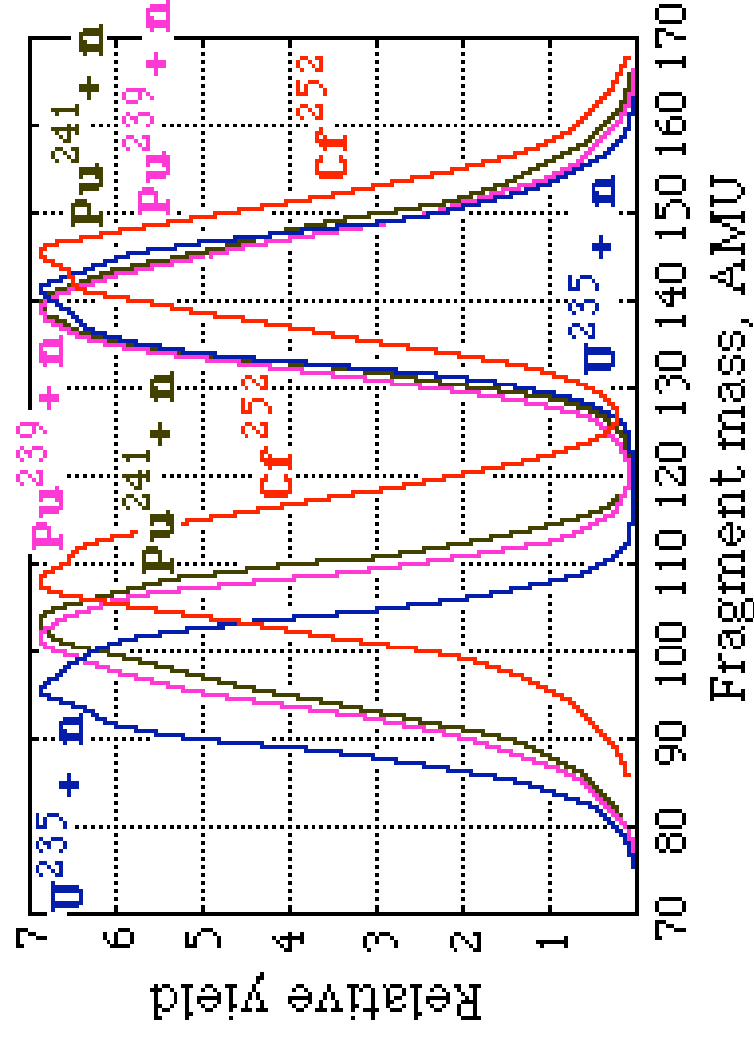
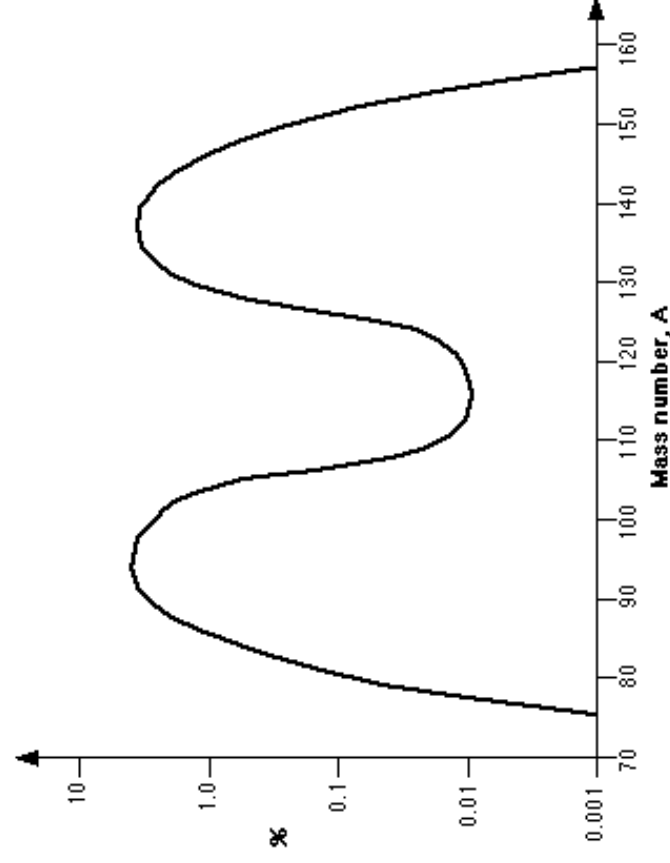
Stabilita je pak dosažena beta-rozpady.

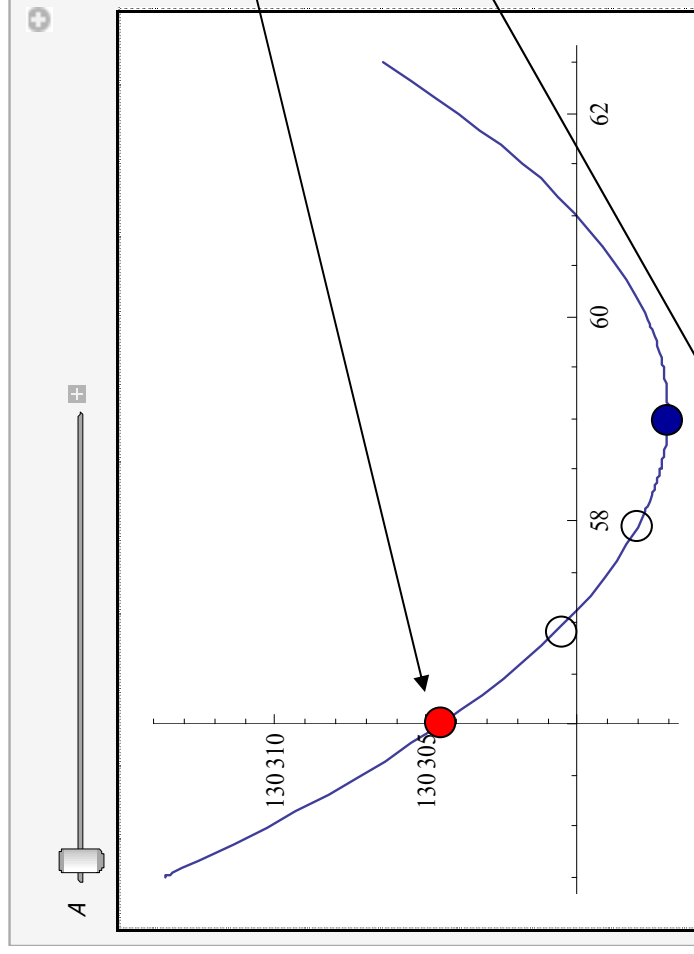
1	H 1																	1B	
2	Li 3	Be 4																	He 2
3	Na 11	Mg 12																	
4	K 19	Ca 20	Sc 21	Ti 22	V 23	Cr 24	Mn 25	Fe 26	Co 27	Ni 28	Cu 29	Zn 30	Ga 31	Ge 32	As 33	Se 34	Br 35	Kr 36	
5	Rb 37	Sr 38	Y 39	Zr 40	Nb 41	Mo 42	Tc 43	Ru 44	Rh 45	Pd 46	Ag 47	Cd 48	In 49	Sn 50	Sb 51	Te 52	I 53	Xe 54	
6	Cs 55	Ba 56	La 57	Hf 72	Ta 73	W 74	Re 75	Os 76	Ir 77	Pt 78	Au 79	Hg 80	Tl 81	Pb 82	Bi 83	Po 84	At 85	Rn 86	
7	Fr 87	Ra 88	Ac 89	Rf 104	Db 105	Sg 106	Bh 107	Hs 108	Mt 109	Ds 110	Rg 111	Uub 112	Uut 113	Uuq 114	Uup 115	Uuh 116	Uus 117	Uuo 118	

Z	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu
	58	59	60	61	62	63	64	65	66	67	68	69	70	71
T	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr
	90	91	92	93	94	95	96	97	98	99	100	101	102	103

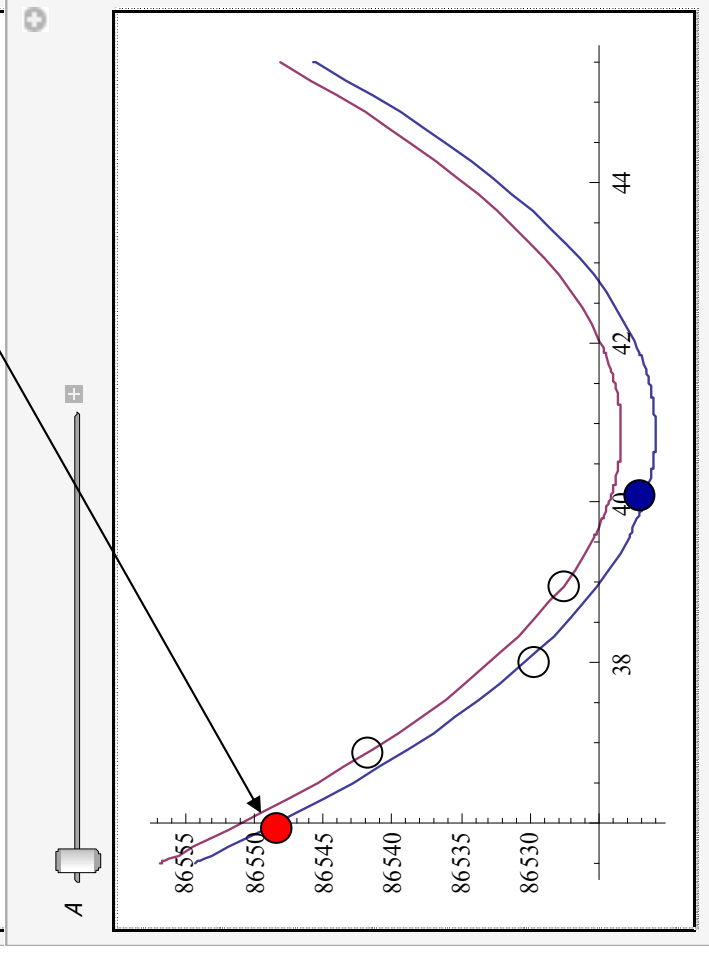
Štěpení je převážně asymetrické, existuje silná tendence produkovat těžké fragmenty s A okolo 140 a 90. Důvodem jsou magická jádra.

Distribution of fission products from Uranium-235





Out[79]=



Out[80]=

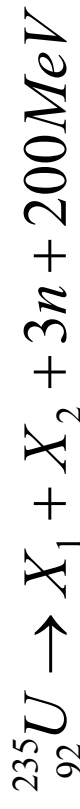
Vzniklá jádra mají příliš mnoho neutronů = nacházejí se mimo údolí stability.
Stabilita je pak dosažena beta-rozpady.

Porovnání chemické a jaderné energie



$$12g \text{ C } 6,023 \cdot 10^{23} \cdot 4,1eV = 6,023 \cdot 10^{23} \cdot 6,56 \cdot 10^{-19} J = 39,5 \cdot 10^4 J =$$

$$395kJ \Rightarrow 1g \text{ C } 32,9kJ$$



$$235g \text{ } {}^{235}_{92}U \text{ } 6,023 \cdot 10^{23} \cdot 200MeV = 6,023 \cdot 10^{23} \cdot 200 \cdot 10^6 \cdot 1,602 \cdot 10^{-19} J =$$

$$1930 \cdot 10^{10} J = 19300GJ \Rightarrow 1g \text{ } {}^{235}_{92}U \text{ } 82,1GJ$$

Deformace jader

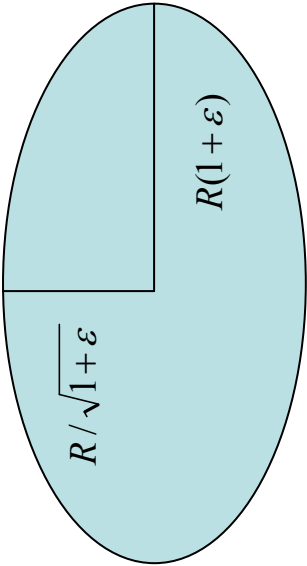
$$V = \frac{4}{3} \pi \left(\frac{R}{\sqrt{1+\varepsilon}} \right)^2 R(1+\varepsilon) = \frac{4}{3} \pi R^3$$

$$\left(\frac{x}{R(1+\varepsilon)} \right)^2 + \left(\frac{y}{R/\sqrt{1+\varepsilon}} \right)^2 = 1 \Rightarrow y(x) = \sqrt{\frac{R^2}{1+\varepsilon} - \frac{x^2}{(1+\varepsilon)^3}} \,;$$

$$S = 4 \pi R^2 (1 + \frac{2}{5} \varepsilon^2)$$

$$R^2 \rightarrow R^2 (1 + \frac{2}{5} \varepsilon^2)$$

$$\frac{1}{R} \rightarrow \frac{1}{R} (1 - \frac{1}{5} \varepsilon^2)$$



```
ln[115]:= "ELIPSOID"

y[x_]:=Sqrt[R^2/(1+ϵ)-x^2/(1+ϵ)^3]

"OBJEM"
Simplify[Series[(y[x])^2,{ϵ,0,4}],R>0&&x>0&&R>x];
(π∫_R^(1+ϵ)_R^(1+ϵ)(%)dx)/(4πR^3/3)

"POVRCH"
Simplify[Series[z[x]Sqrt[1+(∂x y[x])^2},{ϵ,0,4}],R>0&&x>0&&R>x];
(2π∫_R^(1+ϵ)_R^(1+ϵ)(%)dx)/(4πR^2)

Out[115]= ELIPSOID
Out[117]= OBJEM
Out[119]= 1+O(ϵ^5)
Out[120]= POVRCH
Out[122]= 1+2ϵ^2/5-52ϵ^3/105+11ϵ^4/21+O(ϵ^5)
```

$$B(A, Z, \varepsilon) = A \cdot 15,6 \text{ MeV} - A^{2/3} \left(1 + \frac{2}{5} \varepsilon^2 \right) \cdot 17,2 \text{ MeV} - \frac{Z^2}{A^{1/3}} \left(1 - \frac{1}{5} \varepsilon^2 \right) \cdot 0,7 \text{ MeV}$$

$$-\frac{(A-2Z)^2}{A} \cdot 23,3 \text{ MeV} + \begin{matrix} -\frac{1}{A^{1/2}} 12,0 \text{ MeV} & \dots \text{lichá - lichá} \\ 0 \dots \text{lichá - sudá a sudo - lichá} \\ +\frac{1}{A^{1/2}} 12,0 \text{ MeV} & \dots \text{sudá - sudá} \end{matrix}$$

$$B(A, Z, \varepsilon) > B(A, Z, \varepsilon = 0) \Rightarrow \text{Spontánní štěpení}$$

$$-A^{2/3} \left(1 + \frac{2}{5} \varepsilon \right) \cdot 17,2 \text{ MeV} - \frac{Z^2}{A^{1/3}} \left(1 - \frac{1}{5} \varepsilon \right) \cdot 0,7 \text{ MeV} > -A^{2/3} \cdot 17,2 \text{ MeV} - \frac{Z^2}{A^{1/3}} \cdot 0,7 \text{ MeV}$$

$$-A^{2/3} \frac{2}{5} \varepsilon \cdot 17,2 \text{ MeV} + \frac{Z^2}{A^{1/3}} \frac{1}{5} \varepsilon \cdot 0,7 \text{ MeV} > 0$$

$$\frac{Z^2}{A^{1/3}} \frac{1}{5} \cdot 0,7 \text{ MeV} > A^{2/3} \frac{2}{5} \cdot 17,2 \text{ MeV}$$

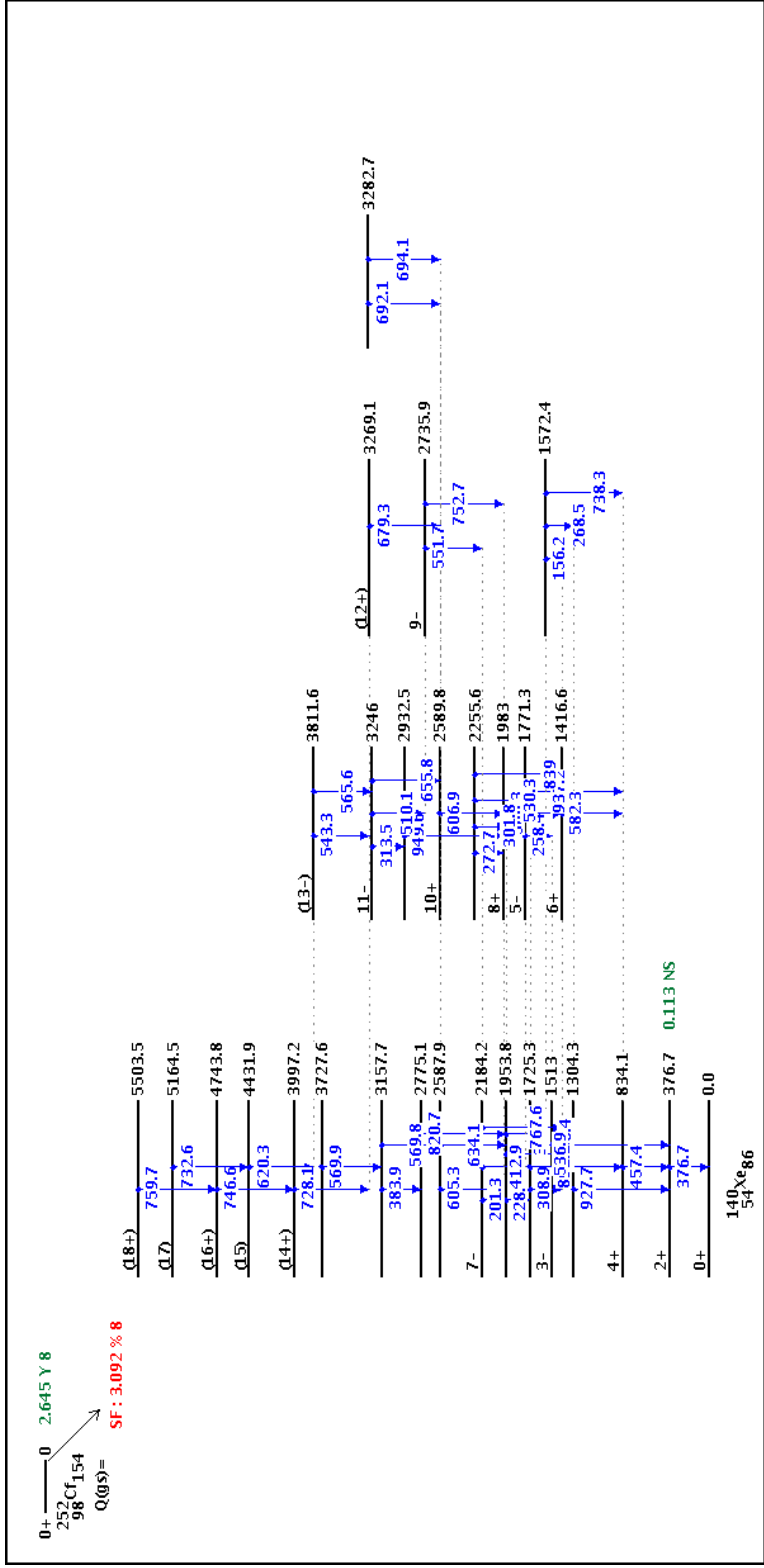
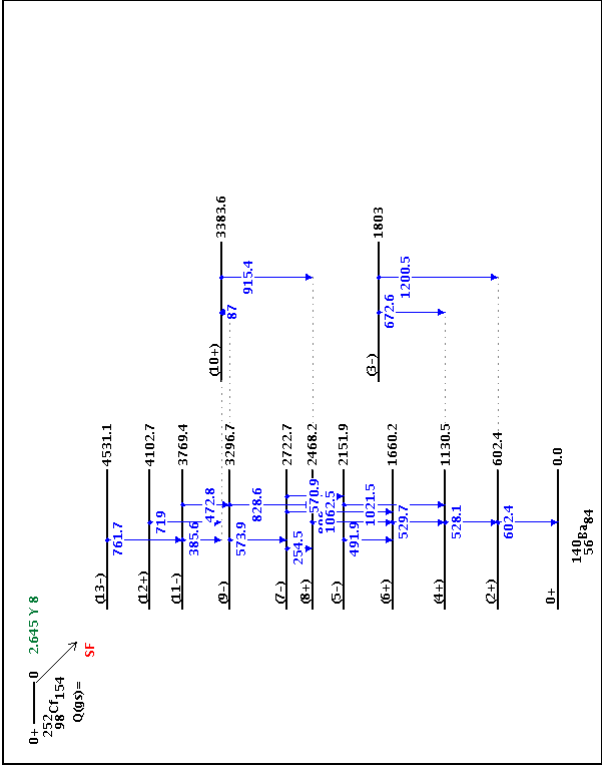
$$\frac{Z^2}{A} > \frac{34,4}{0,7} \approx 49$$

$$^{235}_{92}\text{U} : Z^2 / A = 36$$

$$^{252}_{98}\text{Cf} : Z^2 / A = 38$$

$$^{252}_{98}\text{Cf} \rightarrow ^{140}_{54}\text{Xe} + ^{112-x}_{44}\text{Ru} + x \cdot n$$

$$^{252}_{98}\text{Cf} \rightarrow ^{140}_{56}\text{Ba} + ^{112-x}_{42}\text{Mo} + x \cdot n$$



$$B(A, Z, \varepsilon) = A \cdot 15,6 MeV - A^{2/3} \left(1 + \frac{2}{5} \varepsilon^2 \right) \cdot 17,2 MeV - \frac{Z^2}{A^{1/3}} \left(1 - \frac{1}{5} \varepsilon^2 \right) \cdot 0,7 MeV$$

$$-\frac{(A-2Z)^2}{A} \cdot 23,3 MeV + \frac{1}{A^{1/2}} 12,0 MeV \quad \dots \text{lichó - lichá}$$

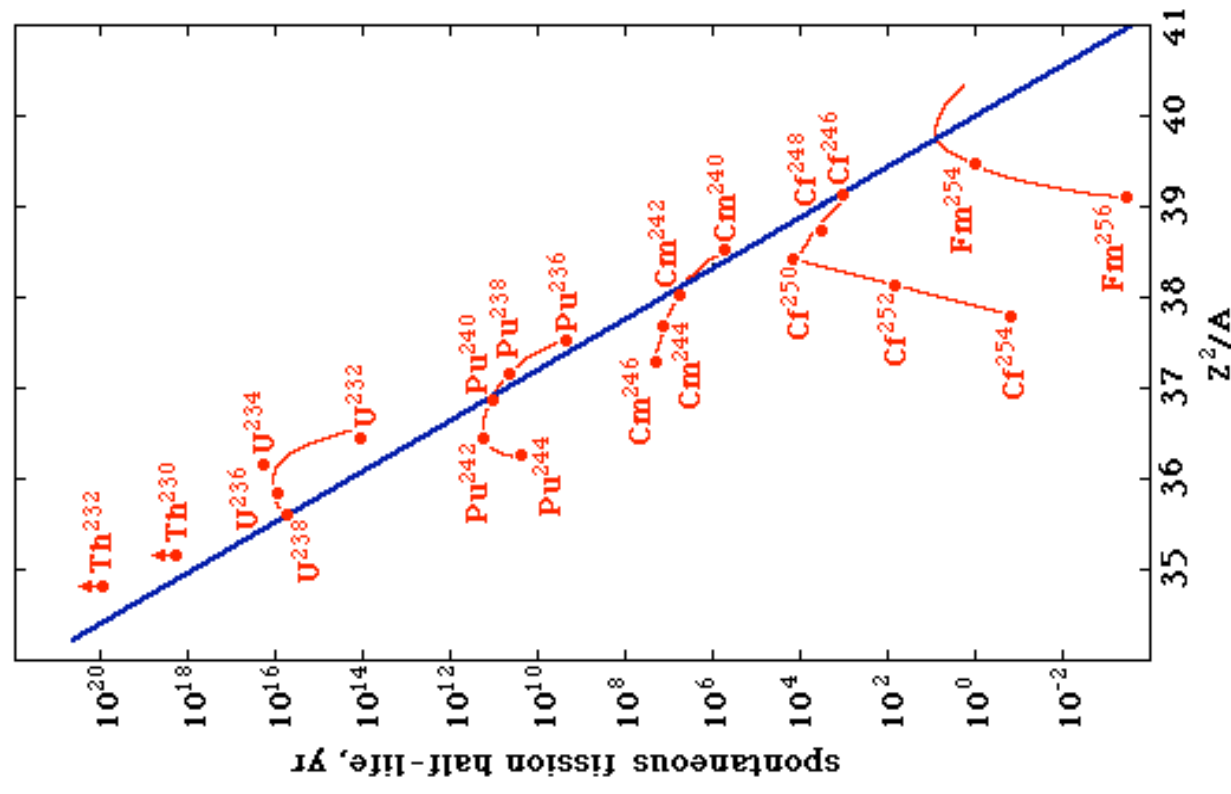
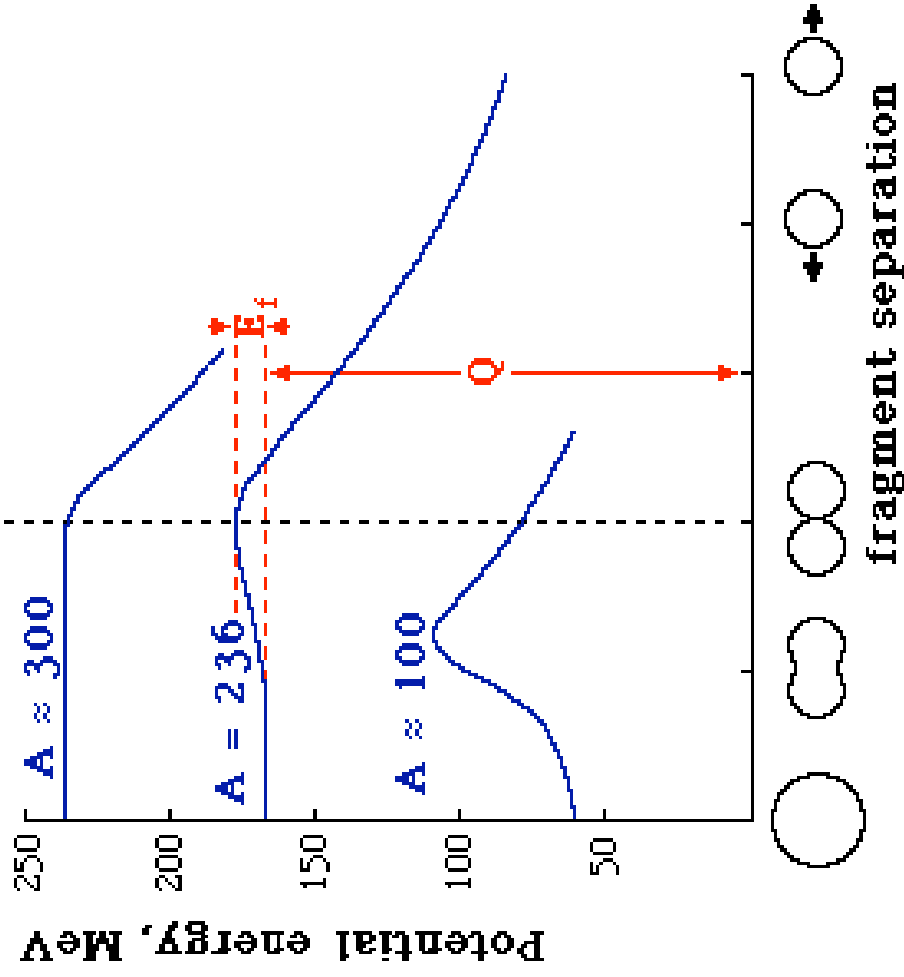
$$+ \frac{1}{A^{1/2}} 12,0 MeV \quad \dots \text{sudá - sudá}$$

$$B(A, Z, \varepsilon) - B(A, Z, \varepsilon = 0) =$$

$$-A^{2/3} \left(1 + \frac{2}{5} \varepsilon^2 \right) \cdot 17,2 MeV - \frac{Z^2}{A^{1/3}} \left(1 - \frac{1}{5} \varepsilon^2 \right) \cdot 0,7 MeV - \left(-A^{2/3} \cdot 17,2 MeV - \frac{Z^2}{A^{1/3}} \cdot 0,7 MeV \right) =$$

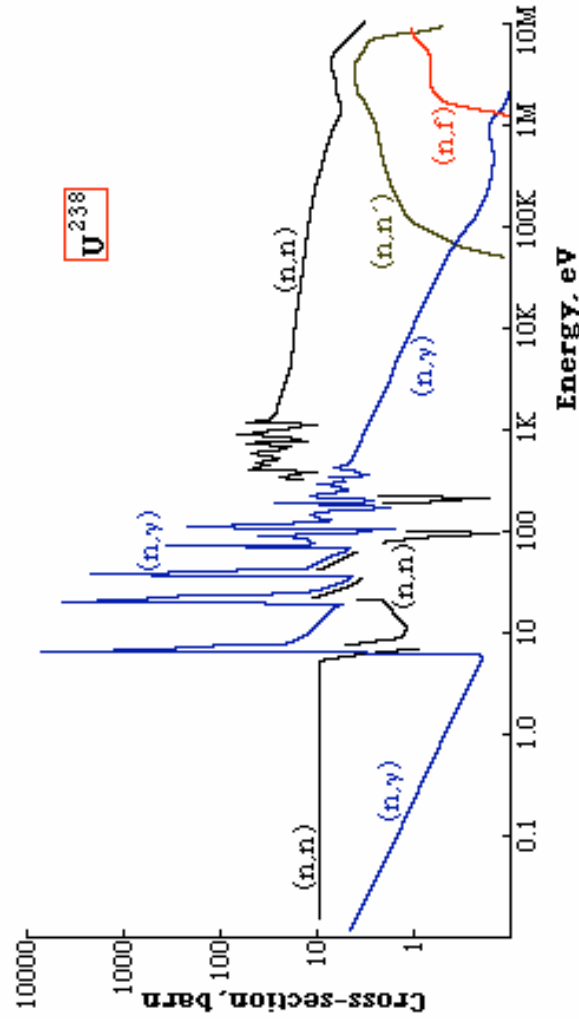
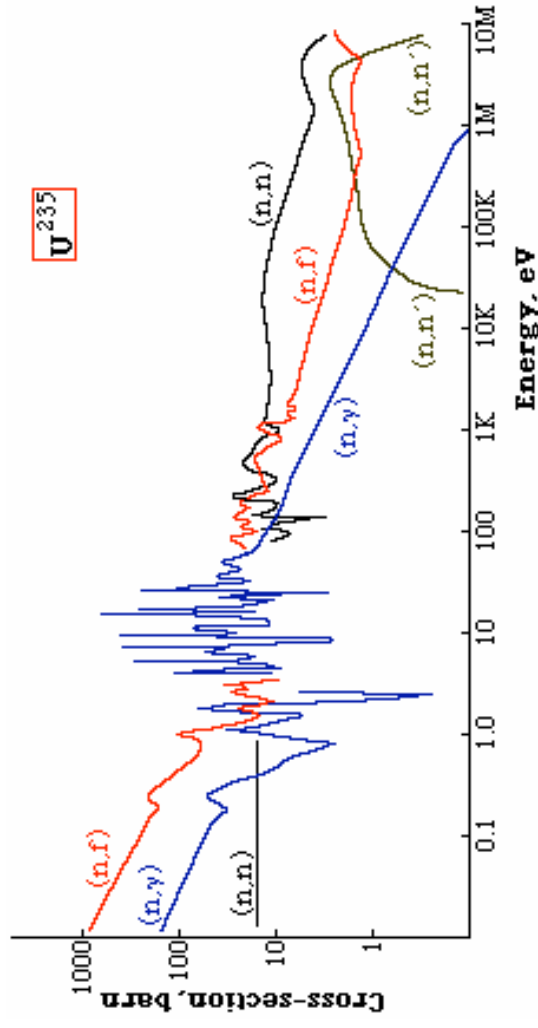
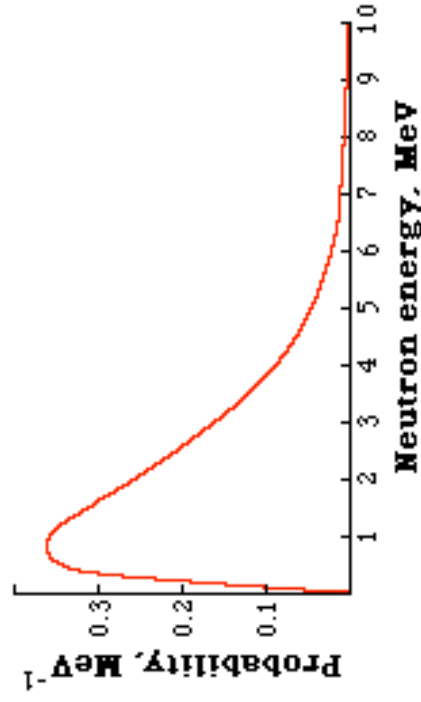
$$-A^{2/3} \frac{2}{5} \varepsilon^2 \cdot 17,2 MeV + \frac{Z^2}{A^{1/3}} \frac{1}{5} \varepsilon^2 \cdot 0,7 MeV$$

$$-235^{2/3} \frac{2}{5} \varepsilon^2 \cdot 17,2 MeV + \frac{92^2}{235^{1/3}} \frac{1}{5} \varepsilon^2 \cdot 0,7 MeV = \varepsilon^2 (-262 + 192) = 70 \varepsilon^2$$



Štěpení U235 lze vzvolat záchytem neutronů

Neutrony je třeba zpomalit (moderovat), aby se zvýšila pravděpodobnost štěpné reakce



Chlazení neutronů = zmenšování jejich kinetické energie

$$n + A \rightarrow n + A; \quad p_n = -p'_n + P_A; \quad \frac{p_n^2}{2m_n} = \frac{p_n'^2}{2m_n} + \frac{P_A^2}{2M_A}$$

$$\left(\frac{\Delta T_k}{T_k} \right)_{\max} = \left(\frac{p_n^2}{2m_n} - \frac{p_n'^2}{2m_n} \right) / \left(\frac{p_n^2}{2m_n} \right) = \left(\frac{P_A^2}{2M_A} \right) / \left(\frac{p_n^2}{2m_n} \right) = \frac{P_A^2}{p_n^2} \frac{m_n}{M_A}$$

$$\frac{p_n^2}{2m_n} = \frac{(P_A - p_n)^2}{2m_n} + \frac{P_A^2}{2M_A} \Rightarrow P_A^2 \left(\frac{1}{2M_A} + \frac{1}{2m_n} \right) = \frac{P_A^2}{p_n^2} \Rightarrow \left(\frac{2M_A}{M_A + m_n} \right)^2$$

$$\left(\frac{\Delta T_k}{T_k} \right)_{\max} = \frac{P_A^2}{p_n^2} \frac{m_n}{M_A} = \left(\frac{2M_A}{M_A + m_n} \right)^2 \frac{m_n}{M_A} \approx \frac{4A}{(A+1)^2}$$

$$\left\langle \frac{\Delta T_k}{T_k} \right\rangle = \frac{1}{2} \left(\frac{\Delta T_k}{T_k} \right)_{\max} = \frac{2A}{(A+1)^2}$$

$$T_{k,0} \rightarrow T_{k,1} = T_{k,0} \left(1 - \frac{2A}{(A+1)^2} \right) \rightarrow T_{k,2} = T_{k,1} \left(1 - \frac{2A}{(A+1)^2} \right) = T_{k,0} \left(1 - \frac{2A}{(A+1)^2} \right)^2 \rightarrow \dots T_{k,N} = T_k \left(1 - \frac{2A}{(A+1)^2} \right)^N$$

$$1 \text{ MeV} \left(1 - \frac{2 \cdot 12}{(12+1)^2} \right)^N = 100 \text{ meV} \Rightarrow (0,858)^N = 10^{-12} \Rightarrow N = \frac{\ln(10^{-12})}{\ln(0,858)} \approx 180$$

$$1 \text{ MeV} \left(1 - \frac{2 \cdot 1}{(1+1)^2} \right)^N = 100 \text{ meV} \Rightarrow (0,5)^N = 10^{-12} \Rightarrow N = \frac{\ln(10^{-12})}{\ln(0,5)} \approx 40$$

Přírodní uran:

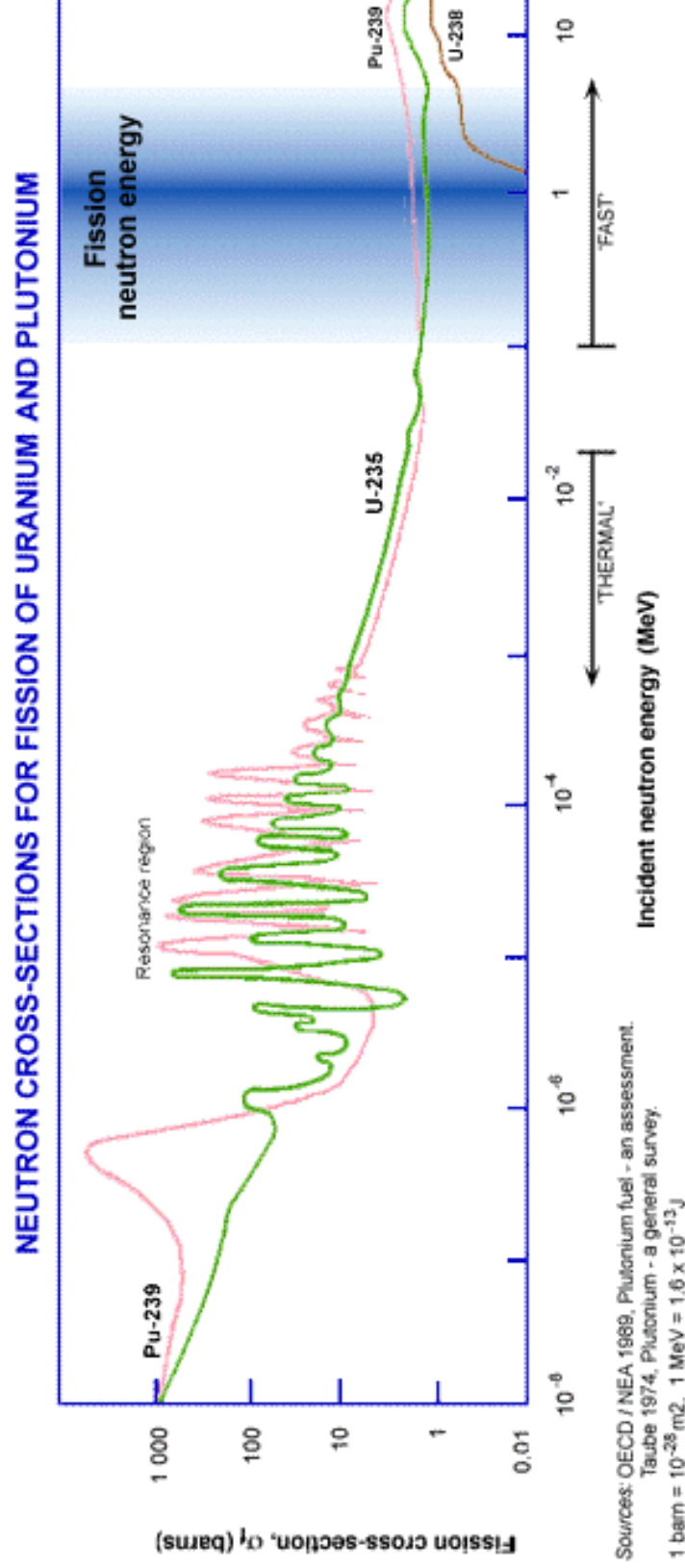
100 neutronů

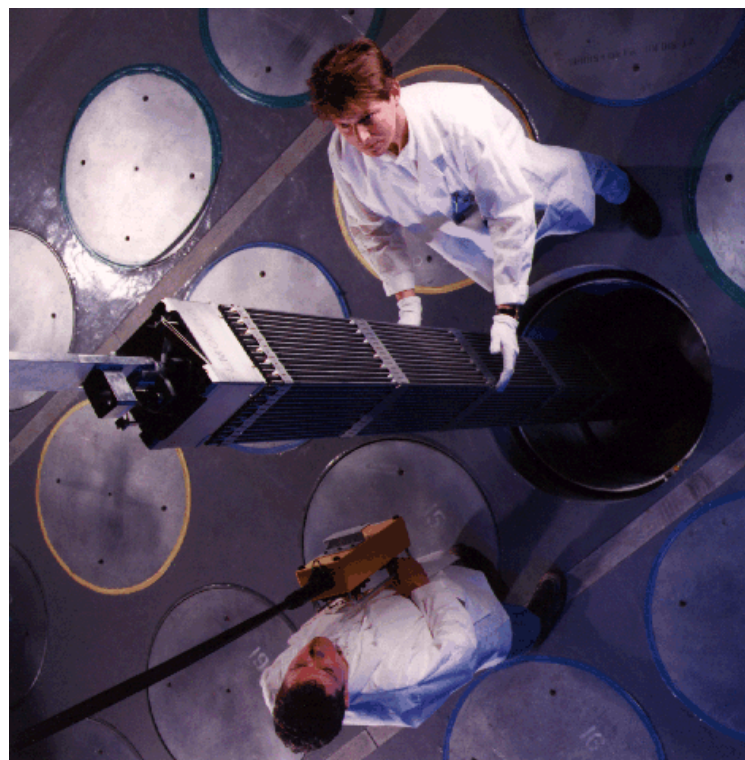
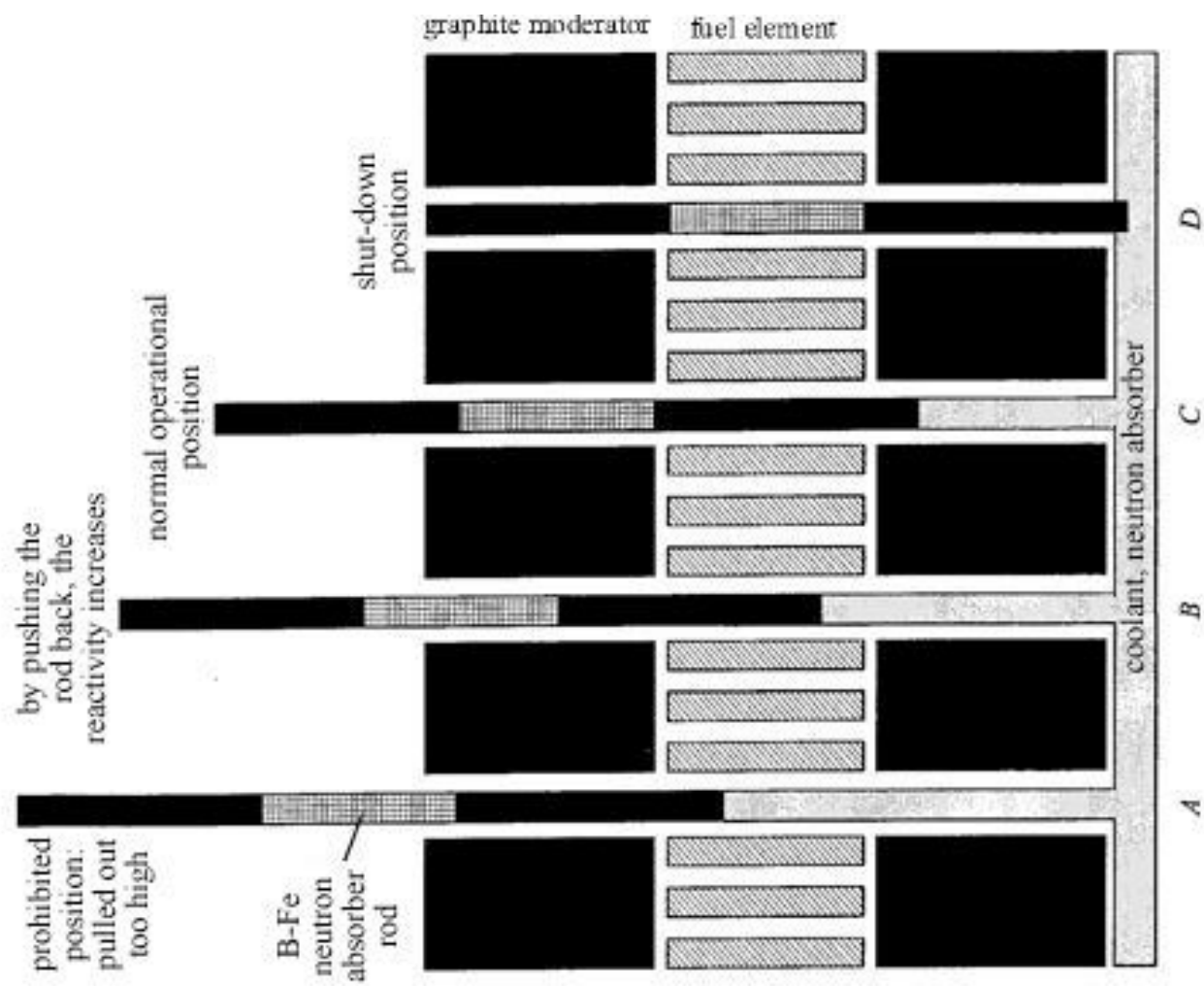
98 je zachyceno U238 a jen 8 z nich vyvolá štěpení x 3 neutrony = 24 neutronů

2 zbývající štěpí U235 x 3 = 6 neutronů

Tj. ze 100 neutronů bude 30

Přírodní uran je nutné obohatit, tj. zvýšit v něm zastoupení U235, přibližně na 4%. To se dělá pomocí např. pomocí odstředivek.





Jaderné reakce



$$Q = M_a + M_A - (M_b + M_B)$$

$$Q > 0$$

Exotermické reakce

$$Q < 0$$

Endotermické reakce



První jaderná reakce objevená Rutherfordem



Chadwickův objev neutronu



Termojaderná fúze – příkladem je Slunce, kde se vodík mění na helium

$$p + p \rightarrow {}^2_1\text{H} + e^+ + \nu_e$$

$$Q = 2m_p - (m_p + m_n - B({}_1^2\text{H}) + m_e) = 2 \cdot 938,27 - (938,27 + 939,57 - 2,2 + 0,511) = 0,4\text{MeV}$$

Neexistuje vázaný stav pp ani nn, pouze pn = deuterium.

$${}^2_1\text{H} + {}^2_1\text{H} \rightarrow {}^4_2\text{He}; Q = 23,8\text{MeV}$$

Ale uvolněná energie je dostatečná na to, aby se přeměnila na odtržení protonu anebo neutronu, takže častěji nastane:

$${}^2_1\text{H} + {}^2_1\text{H} \rightarrow {}^3_2\text{He} + n; Q = 3,3\text{MeV}$$

$${}^2_1\text{H} + {}^2_1\text{H} \rightarrow {}^3_1\text{H} + p; Q = 4,0\text{MeV}$$

Termojaderná fúze – příkladem je Slunce, kde se vodík mění na helium

Po fúzi deuteria by helium mohlo vzniknout přímo v reakci:



Ale uvolněná energie je dostatečná na to, aby se přeměnila na odtržení protonu anebo neutronu, takže častěji nastane:



Další fáze syntézy prvků:

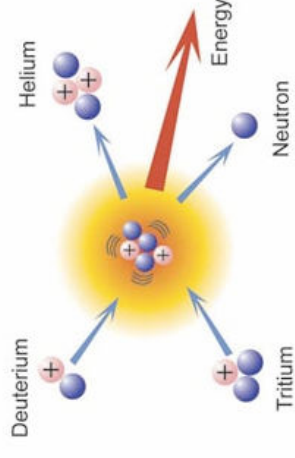
Spalování helia



Další absorpce alfa vedou postupně až ke tvorbě železa

Fúzní reaktory

Využití reakce:



$$P_1(E) \approx e^{-\frac{E}{kT}} \quad \text{Maxwell-Boltzman}$$

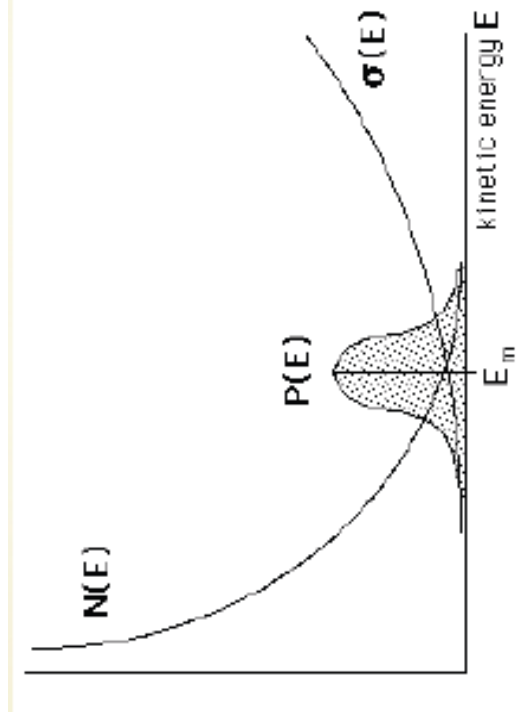
$$P_2(E) \approx e^{-a \int \sqrt{V(x)-E} dx}$$

$$P(E) = P_1(E)P_2(E) \approx e^{-\left(\frac{E}{kT} + a\sqrt{V-E}\right)}$$

$$\frac{d\left(\frac{E}{kT} + a\sqrt{V-E}\right)}{dE} = 0 \Rightarrow \frac{1}{kT} - \frac{a}{2\sqrt{V-E}} = 0$$

Průnik Coulomb. bariery:

$$P_2(E) = e^{-2 \int_R^{R+D} \sqrt{2M(V(x)-E)} \frac{dx}{\hbar c}}$$



$$n_D = n_T = n$$

$$N(t) = n \cdot \sigma \cdot v \cdot t \cdot n$$

$$E(t) = N(t) \cdot Q$$

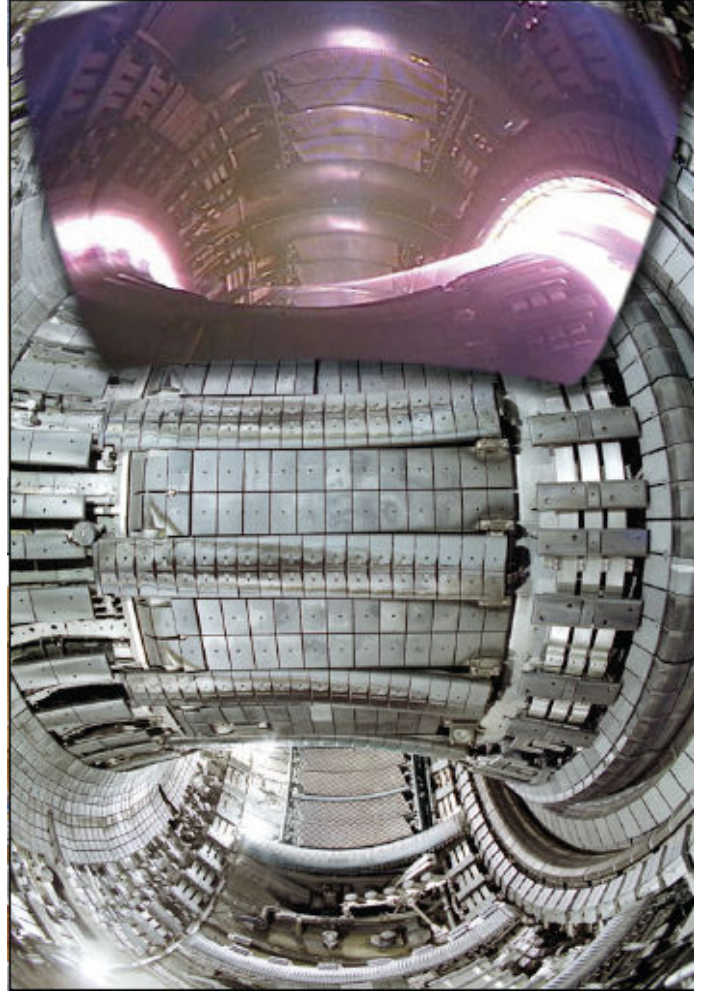
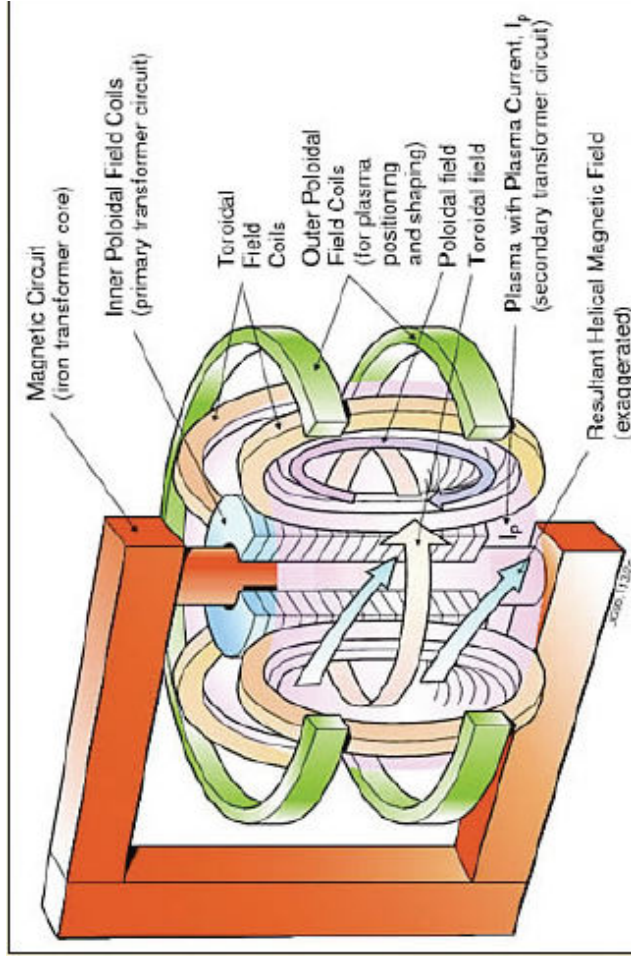
$$E = (n + n) \cdot \frac{3}{2} kT$$

$$E(t) = N(t) \cdot Q > E = (n + n) \cdot \frac{3}{2} kT$$

$$n \cdot \sigma \cdot v \cdot t \cdot n \cdot Q > (n + n) \cdot \frac{3}{2} kT$$

$$n \cdot t > \frac{3kT}{\sigma \cdot v \cdot Q}$$

$$n \cdot t > \frac{3kT}{\sigma \cdot v \cdot Q} = \frac{3 \cdot 25 \cdot 10^{-3} \, eV}{10^{-21} m^3 s^{-1} \cdot 17,6 MeV} \frac{1,2 \cdot 10^8 \, K}{300 K} = \frac{3 \cdot 10 keV}{10^{-21} m^3 s^{-1} \cdot 17,6 MeV} = 1,7 \cdot 10^{20} m^{-3} s^1$$



Objevy nových částic

V roce 1932:

Proton

Neutron

Elektron

Neutrino – postulováno Paulim

Objev pozitronu = Carl D. Anderson



The Nobel Prize in Physics 1936

"for his discovery of
cosmic radiation"



Victor Franz Hess

1/2 of the prize

Austria

Innsbruck University
Innsbruck, Austria

b. 1883
d. 1964

"for his discovery of
the positron"



Carl David Anderson

1/2 of the prize

USA

California Institute of
Technology (Caltech)
Pasadena, CA, USA

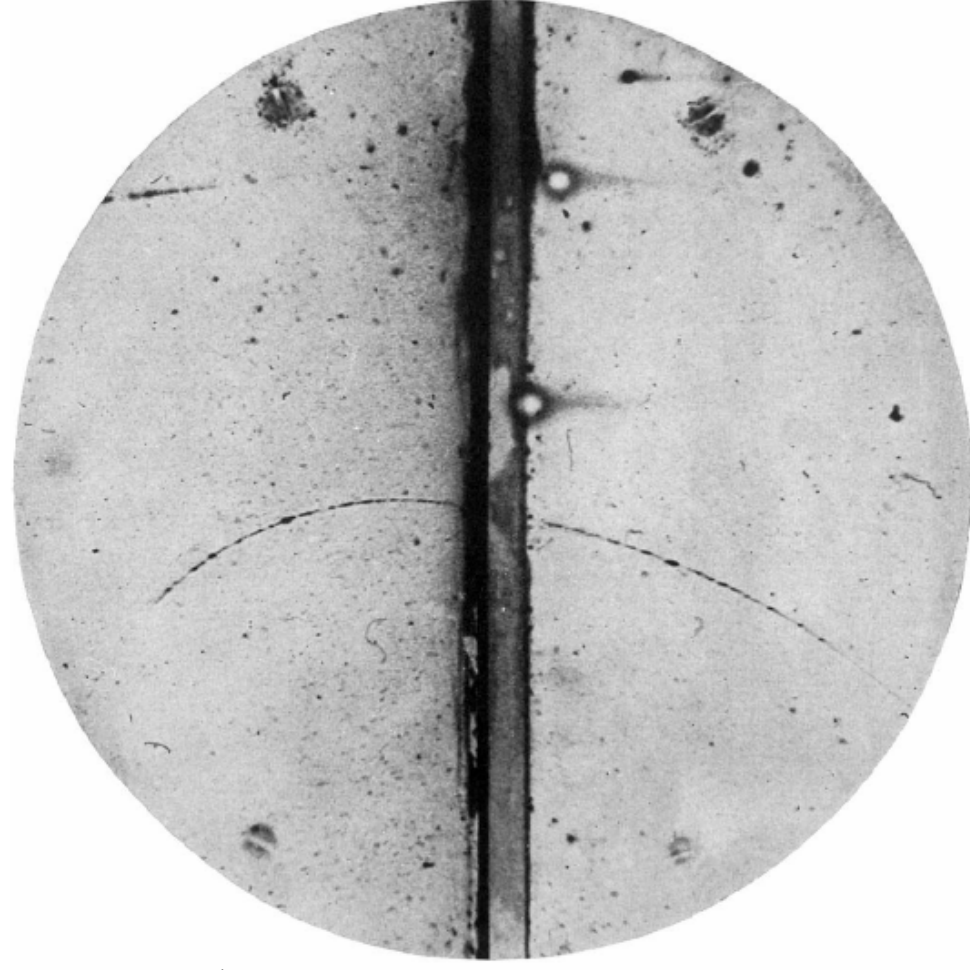
b. 1905
d. 1991

Objev kosmického záření Viktorem Hessem

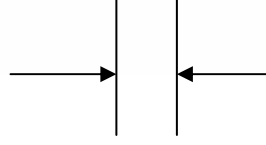
When, in 1912, I was able to demonstrate by means of a series of balloon ascents, that the ionization in a hermetically sealed vessel was reduced with increasing height from the earth (reduction in the effect of radioactive substances in the earth), but that it noticeably increased from 1,000 m onwards, and at 5 km height reached several times the observed value at earth level, I concluded that this ionization might be attributed to the penetration of the earth's atmosphere from outer space by hitherto unknown radiation of exceptionally high penetrating capacity, which was still able to ionize the air at the earth's surface noticeably. Already at that time I sought to clarify the origin of this radiation, for which purpose I undertook a balloon ascent at the time of a nearly complete solar eclipse on the 12th April 1912, and took measurements at heights of two to three kilometres. As I was able to observe no reduction in ionization during the eclipse I decided that, essentially, the sun could not be the source of cosmic rays, at least as far as undeflected rays were concerned.

Objev pozitronu, Carl D. Anderson

Kladně nabitá částice.



P=23 MeV



6 mm Pb

P=63 MeV

FIG. 1. A 63 million volt positron ($H\rho = 2.1 \times 10^5$ gauss-cm) passing through a 6 mm lead plate and emerging as a 23 million volt positron ($H\rho = 7.5 \times 10^4$ gauss-cm). The length of this latter path is at least ten times greater than the possible length of a proton path of this curvature.

$$T_k = \frac{(63MeV)^2}{2 \cdot 938,27MeV} = 2,1MeV;$$

proton: ionizace

$$R_{Pb} = \frac{1}{\rho_{Pb}(dE/dx)_{\min}} \frac{T_k^2}{T_k + m_p} = \frac{1}{11,2 \text{ gcm}^{-3} 1,15MeV/(gcm^{-2})} \frac{2,1^2 MeV^2}{(2,1 + 938,27)MeV}$$

$$= 3,6 \cdot 10^{-4} \text{ cm} = 3,6 \mu m$$

$$T_k = \frac{(23MeV)^2}{2 \cdot 938,27MeV} = 0,28MeV$$

$$R_{Air} = \frac{1}{\rho_{Air}(dE/dx)_{\min}} \frac{T_k^2}{T_k + m_p} = \frac{1}{1,0 \cdot 10^{-3} \text{ gcm}^{-3} 2MeV/(gcm^{-2})} \frac{0,28^2 MeV^2}{(0,28 + 938,27)MeV}$$

$$= 0,042cm = 0,42mm$$

$$\text{Kladný elektron: brzdné záření} \quad E(L) = E_0 e^{-\frac{L}{X_0}} = 63e^{-\frac{6mm}{5,6mm}} = 21,6MeV$$

Částice s neznámou
hmotou m - ionizace

$$\frac{1}{\rho_{Pb}(dE/dx)_{\min}} \frac{(\sqrt{p_1^2 + m^2} - m)^2}{\sqrt{p_1^2 + m^2}} - \frac{1}{\rho_{Pb}(dE/dx)_{\min}} \frac{(\sqrt{p_2^2 + m^2} - m)^2}{\sqrt{p_2^2 + m^2}} = t$$

$$\frac{1}{11,2 \cdot 1,15} \frac{(\sqrt{63^2 + m^2} - m)^2}{\sqrt{63^2 + m^2}} - \frac{1}{11,2 \cdot 1,15} \frac{(\sqrt{23^2 + m^2} - m)^2}{\sqrt{23^2 + m^2}} = 0,6cm$$

$$m = 61,5MeV$$

$$\Delta T \approx \frac{1}{\beta^2}$$

$$\text{In[7]:= sol = solve[1/(11.2*1.15)*((sqrt(63^2+m^2)-m)^2/sqrt(63^2+m^2))-1/(11.2*1.15)*((sqrt(23^2+m^2)-m)^2/sqrt(23^2+m^2))=0.6,m]$$

$$\text{Out[7]= {(m -> -61.5412), (m -> 61.5412)}} \quad \text{sol} = ((23^2 + m^2)/23^2) * (63^2/(63^2 + m^2)) /. sol$$

$$\text{Out[8]= {4.17526, 4.17526}}$$

$$\Delta T_1 / \Delta T_2 = \frac{\beta_2^2}{\beta_1^2} = \frac{\frac{p_2^2}{p_2^2 + m^2}}{\frac{p_1^2}{p_1^2 + m^2}} = \frac{p_2^2(p_1^2 + m^2)}{p_1^2(p_2^2 + m^2)} = \frac{23^2(63^2 + 61,5^2)}{63^2(23^2 + 61,5^2)} \cong 4,1$$