

Oznámení: Proseminář jaderné a subjaderné fyziky,  
každě **Úterý od 9:00 Troja, 9.patro**

## Přednáška 4.

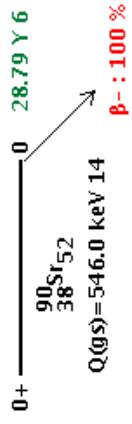
Přirozená radioaktivita, využití pro datování

Kinematika, Mandelstamovy proměnné,

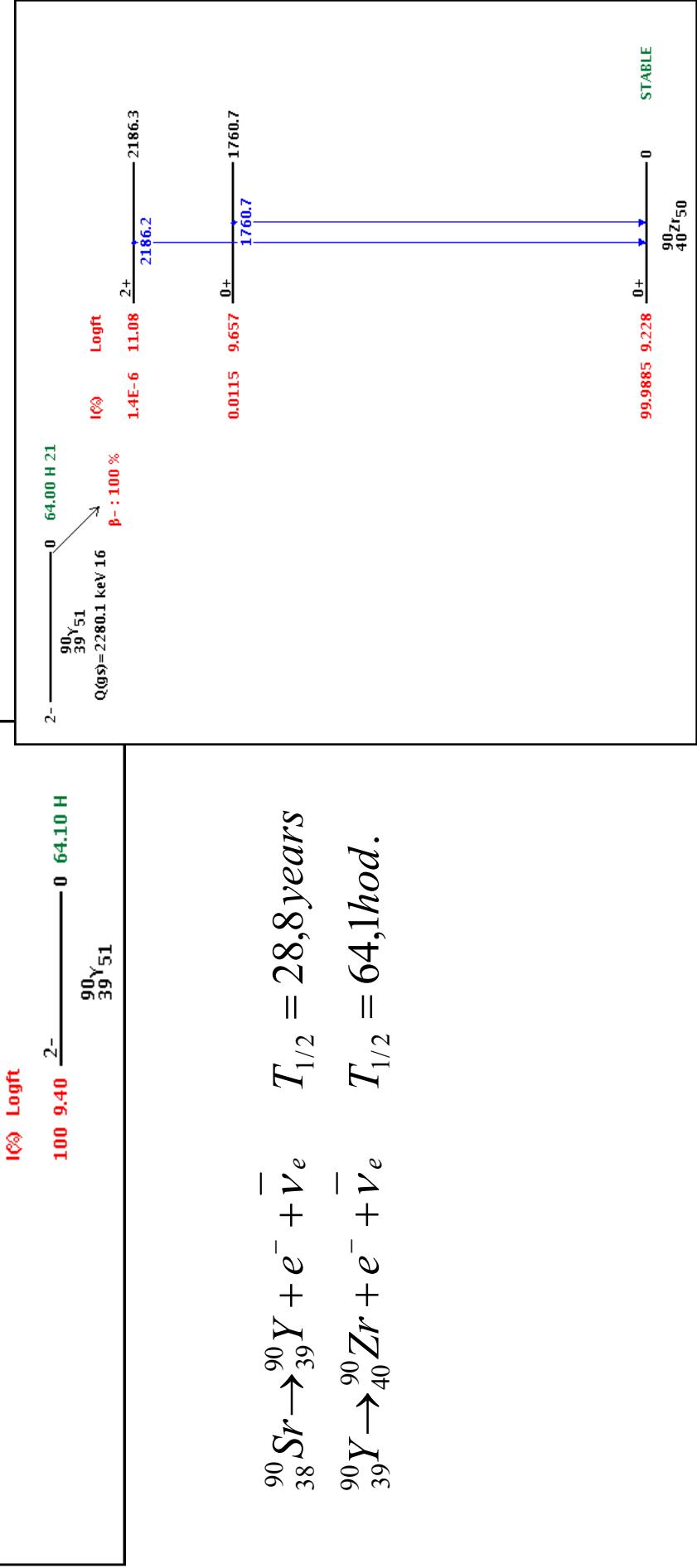
Gama přechody, použití

Interakce jaderného záření s prostředím

Účinky radiace



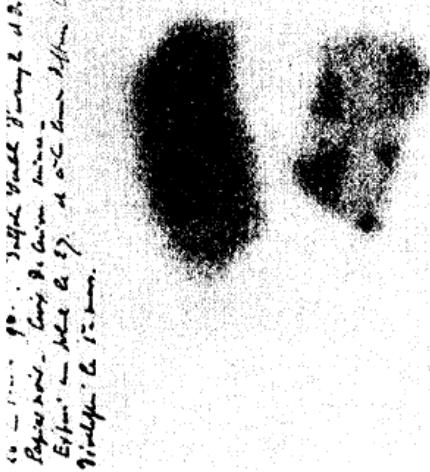
## Příklad beta zářiče



# Objev přirozené radioaktivity:

**Antoine Henri Becquerel** – objevil, že uranové soli produkují pronikavé záření

**Marie a Pierre Curie** zkoumali toto záření a Marie vyřešila záhadu chybějících prvků mezi Bi a U, jako první izolovala Radium a Polonium ve smolinci (pitchblende) z Jáchymova



The Nobel Prize in Physics 1903

"in recognition of the extraordinary services he has rendered by his discovery of spontaneous radioactivity"

"in recognition of the extraordinary services they have rendered by their joint researches on the radiation phenomena discovered by Professor Henri Becquerel"



The Nobel Prize in Chemistry 1911

"in recognition of her services to the advancement of chemistry by the discovery of the elements radium and polonium, by the isolation of radium and the study of the nature and compounds of this remarkable element"



Marie Curie, née Skłodowska

France

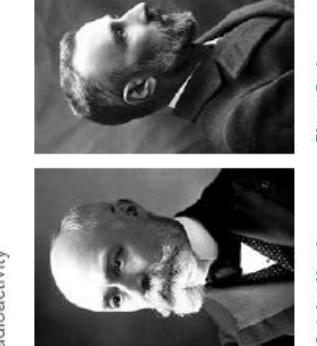
Sorbonne University  
Paris, France

b. 1867  
(in Warsaw, Poland, then  
Russia Empire)  
d. 1934



Pierre Curie  
Marie Curie, née Skłodowska

France



Antoine Henri Becquerel  
Pierre Curie

France

École municipale de physique et de chimie industrielles (Municipal School of Industrial Physics and Chemistry)  
Paris, France

b. 1852  
d. 1908  
b. 1867  
(in Warsaw, Poland, then  
Russia Empire)  
d. 1934

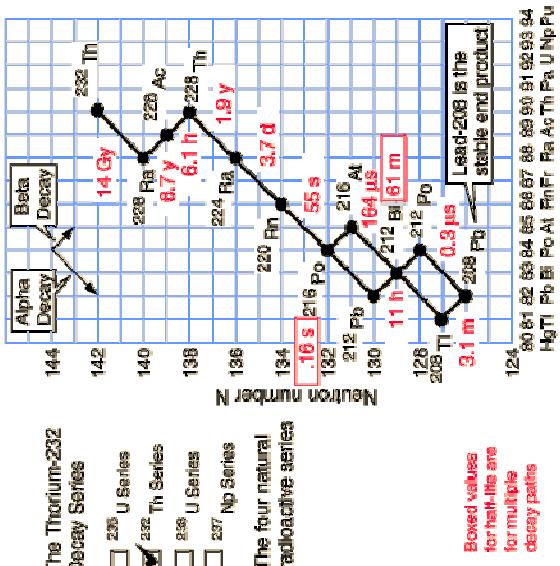
Pierre Curie: „It can even be thought that radium could become very dangerous in criminal hands, and here the question can be raised whether mankind benefits from knowing the secrets of Nature, whether it is ready to profit from it or whether this knowledge will not be harmful for it. The example of the discoveries of Nobel is characteristic, as powerful explosives have enabled man to do wonderful work. They are also a terrible means of destruction in the hands of great criminals who are leading the peoples towards war. I am one of those who believe with Nobel that mankind will derive more good than harm from the new discoveries.“

## Rozpadové řady:

Thoriová:

$$Mod(A,4) = 0:$$

$Mod(A,4) = 0:$	<b>Neptuniová:</b>
$Mod(A,4) = 1:$	<b>Uranová:</b>
$Mod(A,4) = 2:$	<b>Aktiniová:</b>
$Mod(A,4) = 3:$	



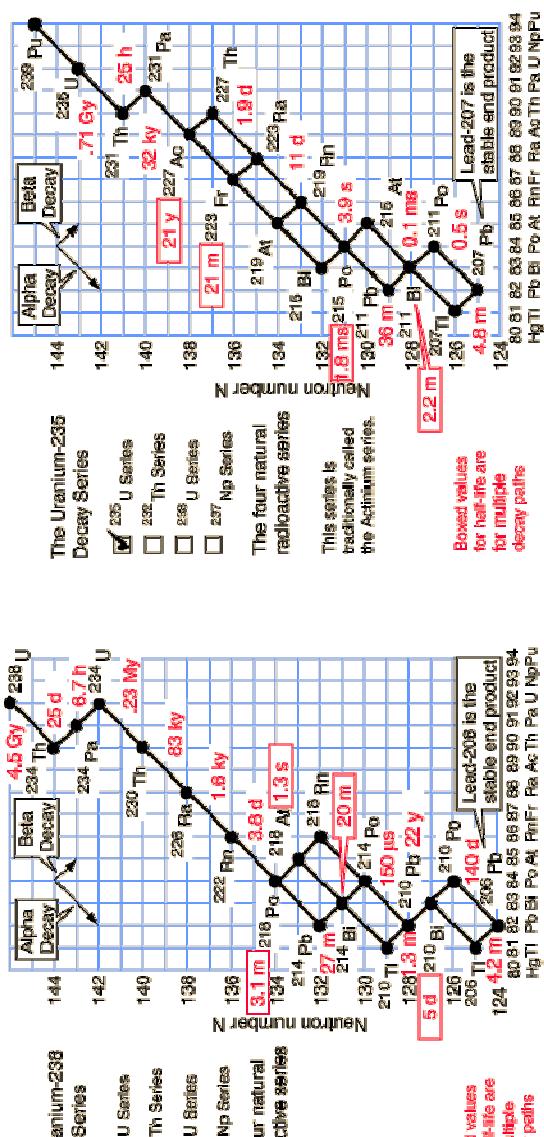
- 235 U Series
- 232 Th Series
- 238 U Series
- 237 Np Series

The four natural radioactive series

$$^{232}_{90}Th \rightarrow \dots$$

Uranová:

$$\begin{array}{l} \text{Mod}(A,4) = 2: \\ \text{Aktinnová:} \\ \text{Mod}(A,4) = 3: \end{array}$$



Boxed values  
for half-life are  
for multiple  
dead paths

Periodická tabulka prvků

Mendeleyev 1869

	1	2	3	4	5	6	7	8
1	H	Li	Be	Mg	Ca	Sr	Ba	Ra
2	He	3	4	12	20	38	56	88
3	Li	Be	3	4	5	6	7	8

18	He	2												
19	H	3	F	9	10	11	12	13	14	15	16	17	18	19
20	B	C	N	O	P	S	Cl	Ar	Kr	Xe	Rn	Fr	Ra	Ac
21	Si	7	8	9	10	11	12	13	14	15	16	17	18	19
22	Al	13	14	15	16	17	18	19	20	21	22	23	24	25
23	Ge	31	32	33	34	35	36	37	38	39	40	41	42	43
24	Ga	31	32	33	34	35	36	37	38	39	40	41	42	43
25	In	49	50	51	52	53	54	55	56	57	58	59	60	61
26	Tl	81	82	83	84	85	86	87	88	89	90	91	92	93
27	Hg	113	114	115	116	117	118	119	120	121	122	123	124	125

6	Ce	58	Pr	60	Nd	61	Pm	62	Sm	63	Eu	64	Tb	65	Dy	66	Ho	67	Er	68	Tm	69	Yb	70	Lu	71		
7	Th	90	Pa	91	U	92	Np	93	Pu	94	Am	95	Bk	96	Cf	97	Es	98	Fm	99	Md	100	No	101	Lr	102	Hg	103

# Datování pomocí radioaktivních rozpadů

## Isotopes of Uranium (Z=92)

Click on an isotope to get more information about its decay

## Zjednodušený příklad:

Isotope	Half-life	Spin	Parity	Decay Mode(s) or Abundance	
<u><a href="#">218U</a></u>	1.5 ms	0+		%A=100	$\frac{t}{T_{235}/\ln(2)}$
<u><a href="#">219U</a></u>	42 us			%A=?	
<u><a href="#">220U</a></u>					
<u><a href="#">221U</a></u>					
<u><a href="#">222U</a></u>	1.0 us	0+		%A=100	$\frac{t}{T_{238}/\ln(2)}$
<u><a href="#">223U</a></u>	18 us			%A=?	
<u><a href="#">224U</a></u>	0.9 ms	0+		%A=100	
<u><a href="#">225U</a></u>	95 ms			%A=100	
<u><a href="#">226U</a></u>	0.35 s	0+		%A=100	
<u><a href="#">227U</a></u>	1.1 m	(3/2+)		%A=100, %EC+%SF < 0.001	$\frac{0.7\%}{99.3\%} = 1 \cdot e^{-\left(\frac{t}{0.70 \cdot 10^9 [\gamma] / \ln(2)} - \frac{1}{4.47 \cdot 10^9 [\gamma] / \ln(2)}\right)}$
<u><a href="#">228U</a></u>	9.1 m	0+		%EC<5, %A>95	
<u><a href="#">229U</a></u>	58 m	(3/2+)		%EC+%SF > 80, %A ~ 20	
<u><a href="#">230U</a></u>	20.8 d	0+		%A=100, %SF < 1.4E-10	
<u><a href="#">231U</a></u>	4.2 d	(5/2-)		%EC=100, %A ~ 0.0055	
<u><a href="#">232U</a></u>	68.9 y	0+		%A=100, %24NE=9E-11 7	
<u><a href="#">233U</a></u>	1.592e+5 y	5/2+		%A=100, %SF < 6E-11, %24NE < 9.5E-11	
<u><a href="#">234U</a></u>	2.455e+5 y	0+		%Abundance=0.0055 5, %A=100, %SF=1.64E-9 22, % NE=9E-12 7, %MG=1.4E-11 3	
<u><a href="#">235U</a></u>	7.038e+8 y	7/2-		%Abundance=0.7200 12, %A=100, %SF=7.0E-9 21, % 20NE=8E-10 4	
<u><a href="#">236U</a></u>	25 m	1/2+		%IT=100	<u><a href="#">238mU</a></u> 1 ns
<u><a href="#">236U</a></u>	2.342e7 y	0+		%A=100, %SF=9.4E-8 4	<u><a href="#">239U</a></u> 23.45 m 5/2+
<u><a href="#">237U</a></u>	6.75 d	1/2+		%B=100	<u><a href="#">240U</a></u> 14.1 h 0+
<u><a href="#">238U</a></u>	4.468e+9 y	0+		%Abundance=99.2745 60, %A=100, %SF=5.45E-5 7, % BB=2.2E-10 7	<u><a href="#">241U</a></u>
<u><a href="#">238mU</a></u>	225 ns	0+			<u><a href="#">242U</a></u> 16.8 m 0+

## Uran.nb

# Určování stáří hornin

## Metoda Uran-Olovo

$^{238}\text{U}$	miliardy let	4.468	$^{206}\text{Pb}$
$^{235}\text{U}$		0.704	$^{207}\text{Pb}$
$^{40}\text{K}$		1.251	$^{40}\text{Ar}$
$^{87}\text{Rb}$		48.8	$^{87}\text{Sr}$

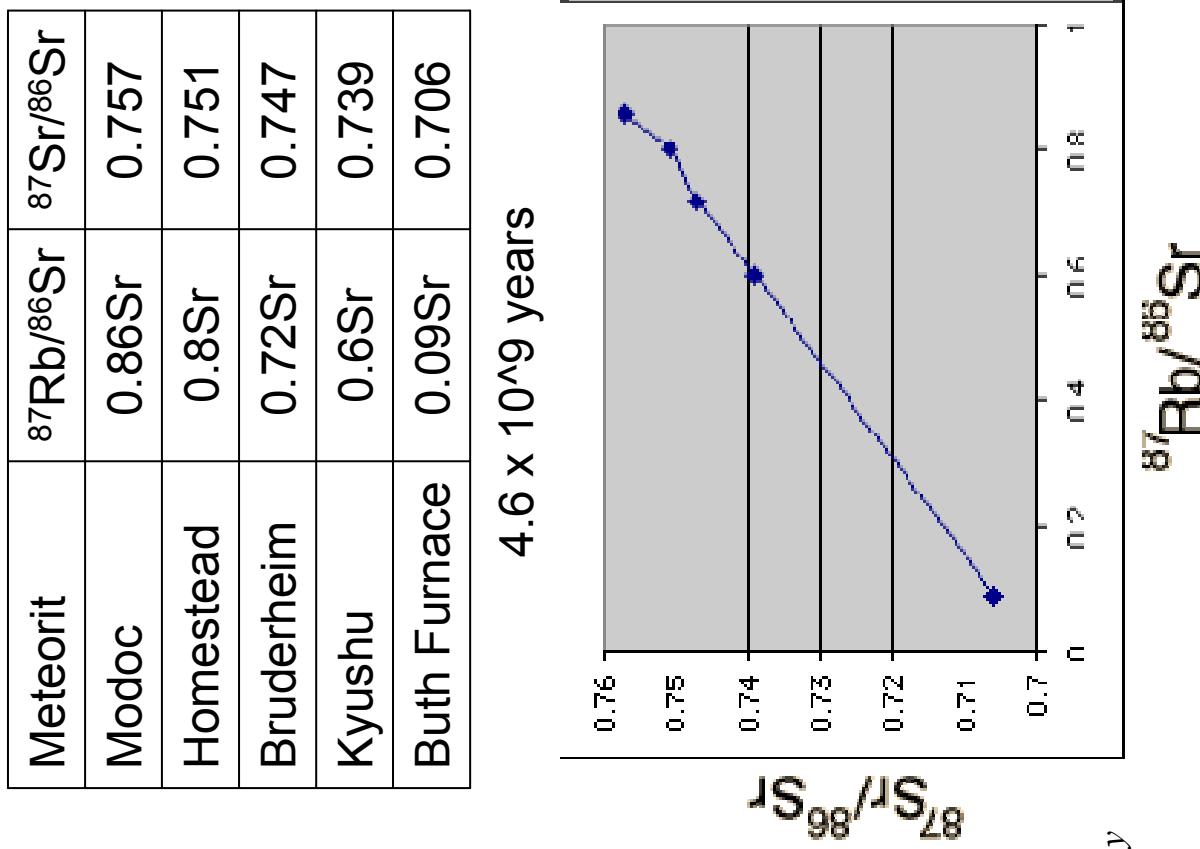
-V uranových dolech: olovo vzniklo z uranu, nebožlo tam žádné na začátku, protože se chemicky nesnese s uranem  
-Stáří Země určeno jako 4,2 miliardy let.

## Metoda Rubidium-Stroncium

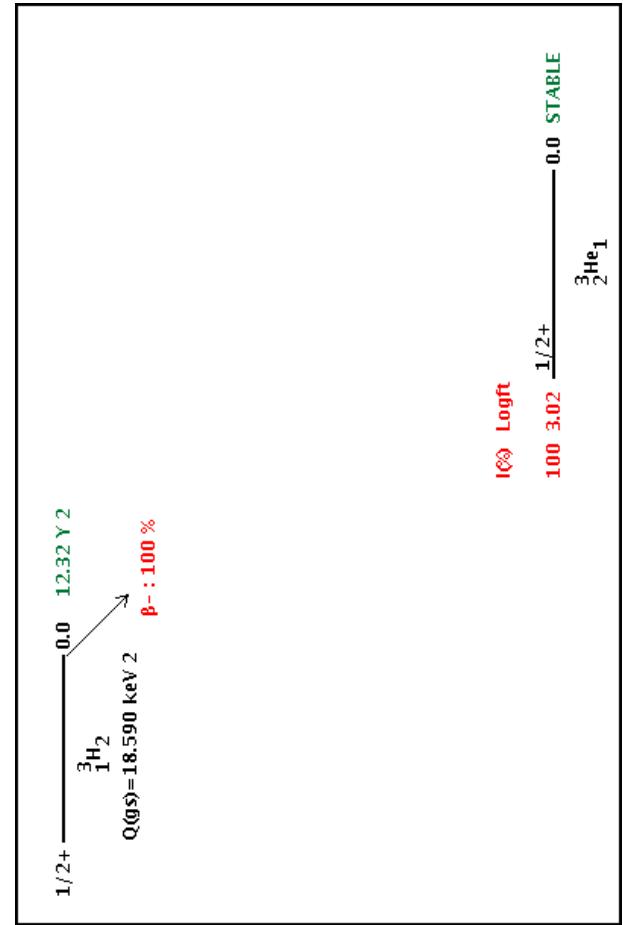
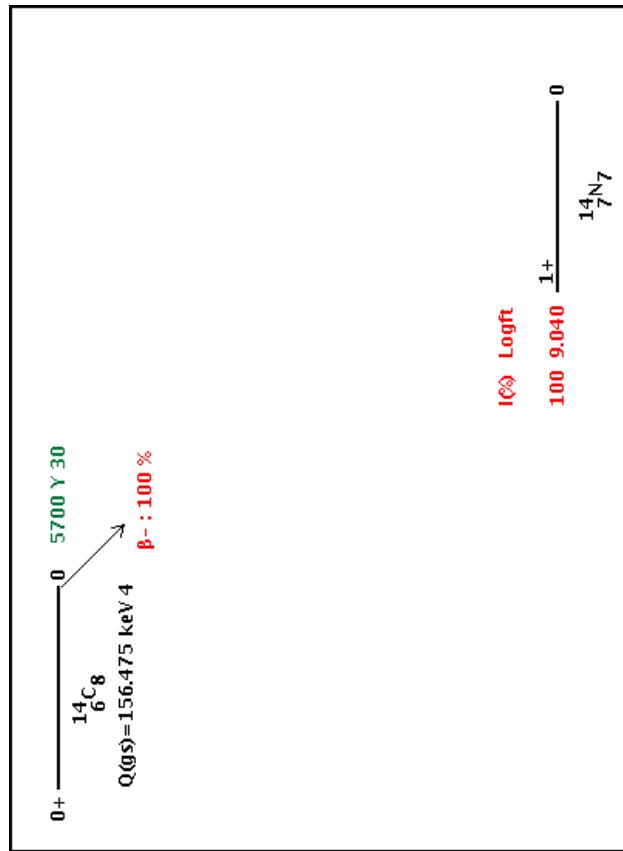
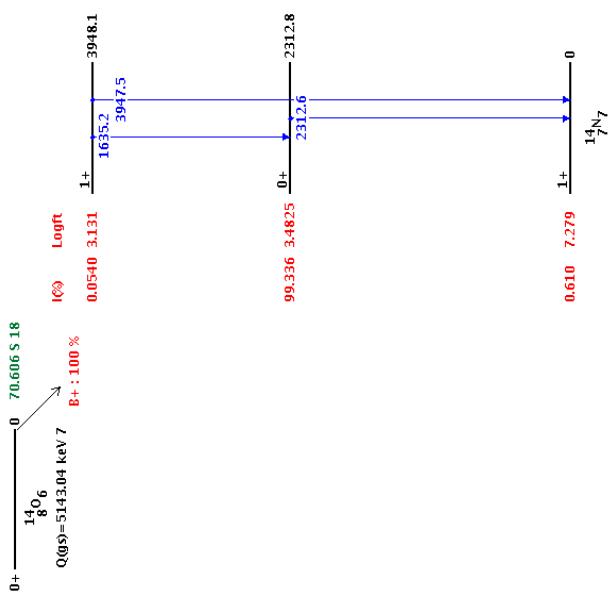
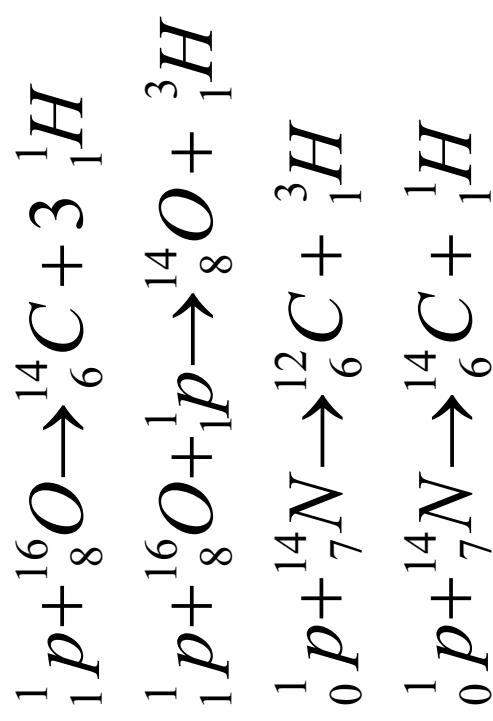
# Metoda Rubidium-Stroncium

$$\begin{aligned} {}_{38}^{86}Sr & \\ {}_{37}^{87}Rb \rightarrow {}_{38}^{87}Sr + e^- + \bar{\nu}_e & \quad T_{1/2} = 46,8 \cdot 10^9 \text{ years} \end{aligned}$$

$$\begin{aligned} t = 0: & \left( \frac{{}_{38}^{87}Sr}{{}_{38}^{86}Sr} \right)_{t=0}, \left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_{t=0} \\ \left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_t &= \frac{\left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_{t=0} \cdot e^{-\frac{t}{T_{1/2}/\ln(2)}}}{\left( \frac{{}_{38}^{86}Sr}{{}_{38}^{86}Sr} \right)_{t=0}} = \left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_{t=0} \cdot e^{-\frac{t}{T_{1/2}/\ln(2)}} \\ \left( \frac{{}_{38}^{87}Sr}{{}_{38}^{86}Sr} \right)_t &= \frac{\left( \frac{{}_{38}^{87}Sr}{{}_{38}^{86}Sr} \right)_{t=0} + \left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_{t=0} \cdot \left( 1 - e^{-\frac{t}{T_{1/2}/\ln(2)}} \right)}{\left( \frac{{}_{38}^{86}Sr}{{}_{38}^{86}Sr} \right)_{t=0}} = \\ \left( \frac{{}_{38}^{87}Sr}{{}_{38}^{86}Sr} \right)_t &= \frac{\left( \frac{{}_{38}^{87}Sr}{{}_{38}^{86}Sr} \right)_{t=0} + \left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_{t=0} \cdot \left( e^{\frac{t}{T_{1/2}/\ln(2)}} - 1 \right)}{\left( \frac{{}_{38}^{86}Sr}{{}_{38}^{86}Sr} \right)_{t=0}} = \\ \left( \frac{{}_{38}^{87}Sr}{{}_{38}^{86}Sr} \right)_t &+ \left( \frac{{}_{37}^{87}Rb}{{}_{38}^{86}Sr} \right)_{t=0} \cdot \left( e^{\frac{t}{T_{1/2}/\ln(2)}} - 1 \right) \end{aligned}$$

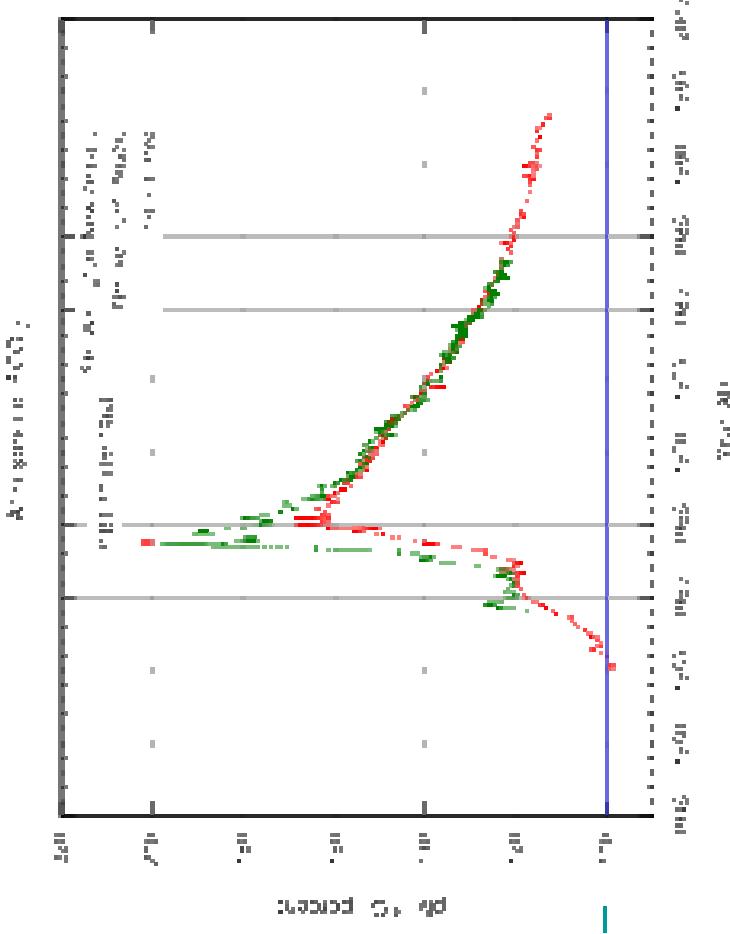


## Radioaktivní izotopy vznikající v atmosféře působením kosmického záření:



# Datování stáří živých organismů pomocí uhlíku 14

- Uhlík 14 tvoří CO<sub>2</sub> a ten se dostává do zelených rostlin
- Živočichové, kteří se žíví zelenými rostlinami a-nebo jinými živočichy, kteří pojídají zelenou stravu dostávají do těla C14, ten je v rovnováze s C12
- Jakmile umřou, začne se obsah C14 v jejich tělech snižovat



Je třeba brát do úvahy změny obsahu uhlíku 14 v atmosféře



Atmospheric 14C, New Zealand[1] and Austria[2]. The New Zealand curve is representative for the Southern Hemisphere, the Austrian curve is representative for the Northern Hemisphere. Atmospheric nuclear weapon tests almost doubled the concentration of <sup>14</sup>C in the Northern Hemisphere [3].

# Kinematika – 4 vektory

4 – vektor

$$\begin{aligned}(\vec{V}_0, \vec{V}) &= (V_0, V_x, V_y, V_z) \\V_0 &\rightarrow \gamma(V_0 - \beta V_x) \\ \vec{V}\end{aligned}$$

$$V_x \rightarrow \gamma(V_x - \beta V_0)$$

Součin dvou 4-vektorů:

$$\begin{aligned}P \cdot Q &= P_0 Q_0 - \vec{P} \cdot \vec{Q} && \text{Součin dvou 4-vektorů je invariantní při Lorentzově transformaci:} \\P \cdot Q &\rightarrow \gamma(P_0 - \beta P_x)\gamma(Q_0 - \beta Q_x) - \gamma(P_x - \beta P_0)\gamma(Q_x - \beta Q_0) - P_y Q_y - P_z Q_z = \\P_0 Q_0 (\gamma^2 - \gamma^2 \beta^2) + P_0 Q_x (-\gamma^2 \beta + \gamma^2 \beta) + P_x Q_0 (-\gamma^2 \beta^2 + \gamma^2 \beta^2) - P_x Q_x (-\gamma^2 \beta^2 + \gamma^2 \beta^2) - P_y Q_y - P_z Q_z = \\P_0 Q_0 (\gamma^2 - \gamma^2 \beta^2) - P_x Q_x (\gamma^2 - \gamma^2 \beta^2) - P_y Q_y - P_z Q_z &= P_0 Q_0 - P_x Q_x - P_y Q_y - P_z Q_z = P \cdot Q\end{aligned}$$

Příkladem 4-vektoru je:

$$(ct, r)$$

**4-vektor energie a hybnosti, 4-impuls:**

$$P = (E, \vec{p})$$

$$P \cdot P = E \cdot E - \vec{p} \cdot \vec{p} = E^2 - p^2 = M^2$$

## Použití 4-hybnosti pro výpočet prahu reakcí. Příklad - objev antiprotonu

$$p + p \rightarrow p + p + p + \bar{p}$$

$$(E, \vec{p}) + (m_p, \vec{0}) = (\sum_{i=1}^4 E_i, \sum_{i=1}^4 \vec{p}_i)$$

$$P_{init} = (E + m_p, \vec{p})$$

$$P_{init}^2 = (E + m_p)^2 - |\vec{p}|^2 = E^2 - |\vec{p}|^2 + m_p^2 + 2m_p E = 2m_p^2 + 2m_p E$$

$$P_{thr} = (\sum_{i=1}^4 m_i, \vec{0}) = (4m_p, \vec{0})$$

$$P_{thr}^2 = 16m_p^2$$

$$P_{init}^2 = P_{thr}^2$$

$$2m_p^2 + 2m_p E_{thr} = 16m_p^2$$

$$E_{thr} = 7m_p$$

$$T_{k,thr} = E_{thr} - m_p = 6m_p$$

# Kinematika - Mandelstamovy invarianty s a t

$$s = (P_1 + P_2)^2$$

*cms:*

$$P_1 = (E_1, \vec{p}^*), P_2 = (E_2, \vec{p}^*)$$

$$s = (E_1 + E_2)^2 - (\vec{p}^* - \vec{p}^*)^2 = (E_1 + E_2)^2 = E_{cms}^2 > 0$$

$$t = (P_3 - P_1)^2$$

*cms:*

$$P_1 = (E_1, \vec{p}_1^*), P_3 = (E_3, \vec{p}_3^*)$$

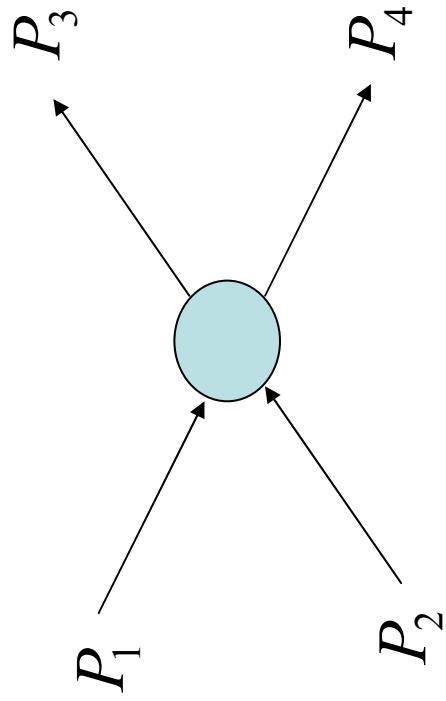
$$t = (E_1 - E_3)^2 - (\vec{p}_1^* - \vec{p}_3^*)^2 = M_1^2 + M_3^2 - 2(E_1 E_3 - p_1^* p_3^*)$$

$$M_1 = M_3 = M \Rightarrow E_1 = E_3 = E;$$

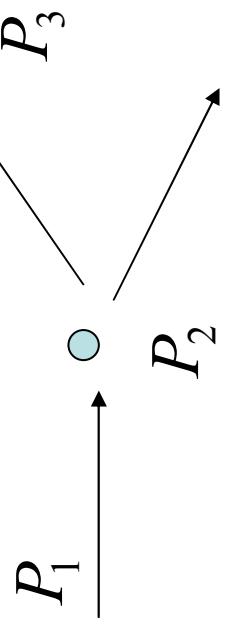
$$M_1 = M_3 \& M_2 = M_4 \Rightarrow \left| \vec{p}_1^* \right| = \left| \vec{p}_3^* \right| = \left| \vec{p}^* \right|$$

$$t = 2M^2 - 2 \left( E^2 - \left| \vec{p}^* \right|^2 \cos(\theta^*) \right) = -2 \left( (E^2 - M^2) - \left| \vec{p}^* \right|^2 \cos(\theta^*) \right) =$$

$$-2 \left| \vec{p}^* \right|^2 (1 - \cos(\theta_{13})) = -4 \left| \vec{p}^* \right|^2 \sin^2(\theta^*/2) = -\left| \vec{q} \right|^2$$



## Kinematika - Mandelstamovy invarianty s a t



$$t = (P_3 - P_1)^2 = (P_4 - P_2)^2$$

lab:

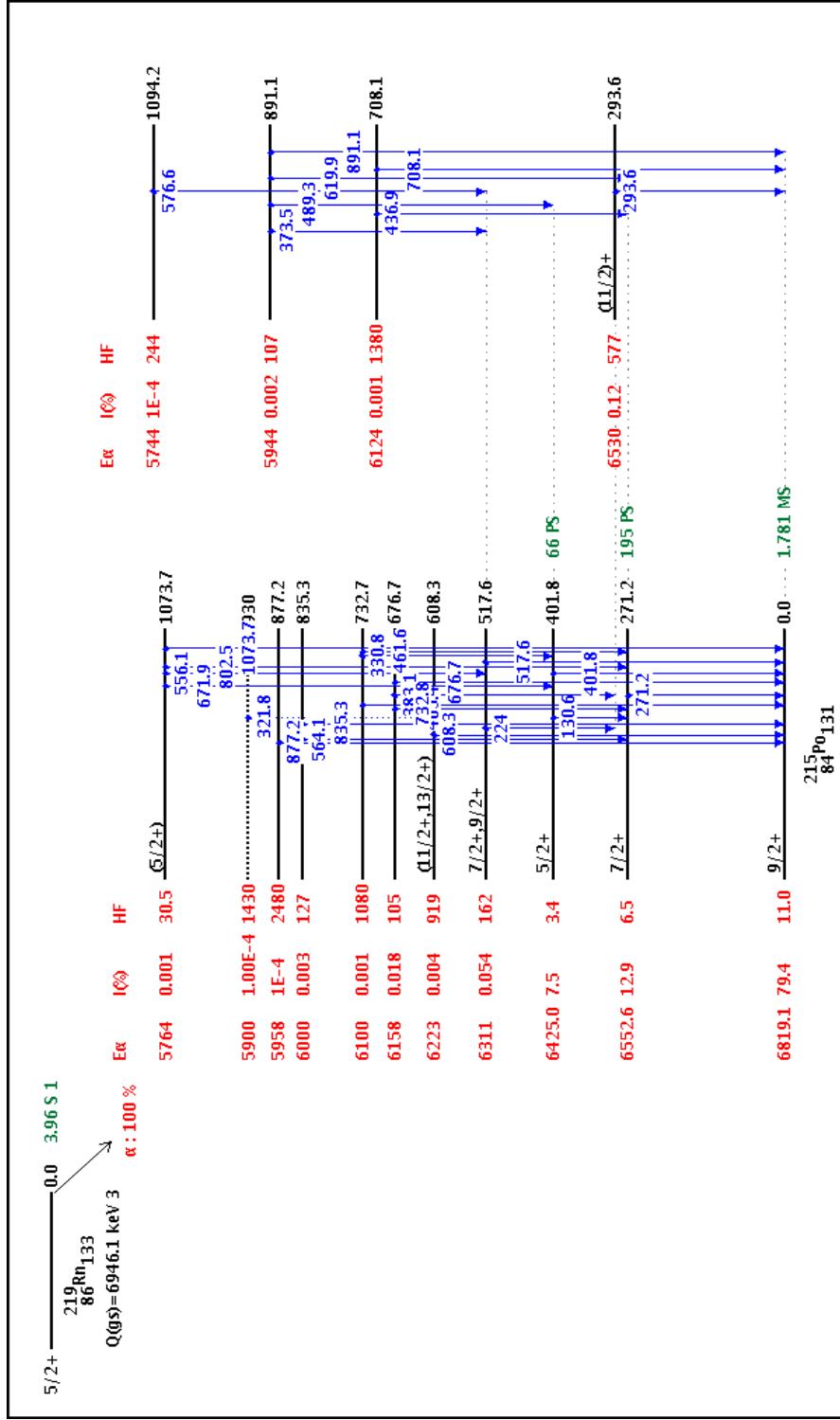
$$P_4 = (E, \vec{p}), P_2 = (m, \vec{0})$$

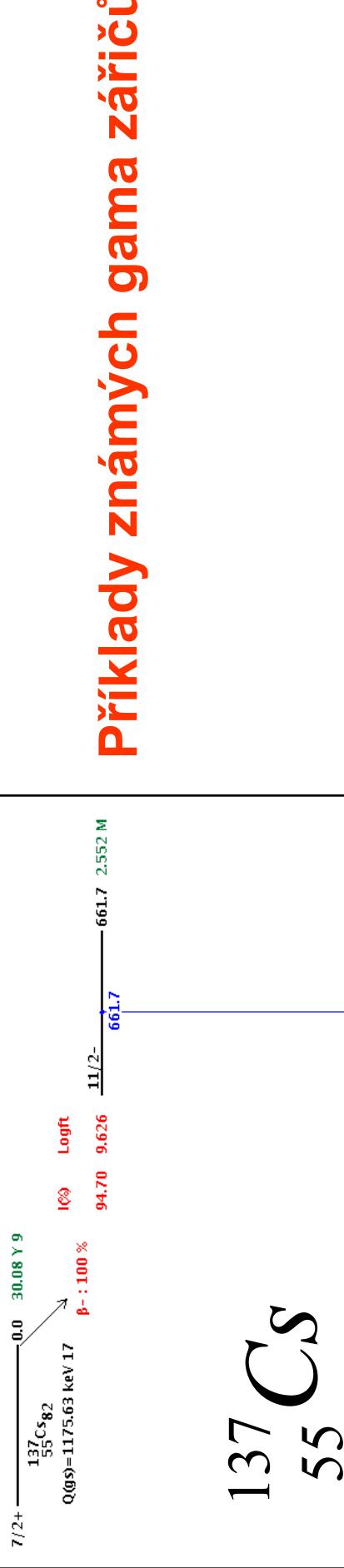
$$t = (E - m)^2 - (\vec{p})^2 = m^2 - 2Em + m^2 = -2m(E - m) = -2mT_k$$

$$T_k = \frac{-t}{2m}$$

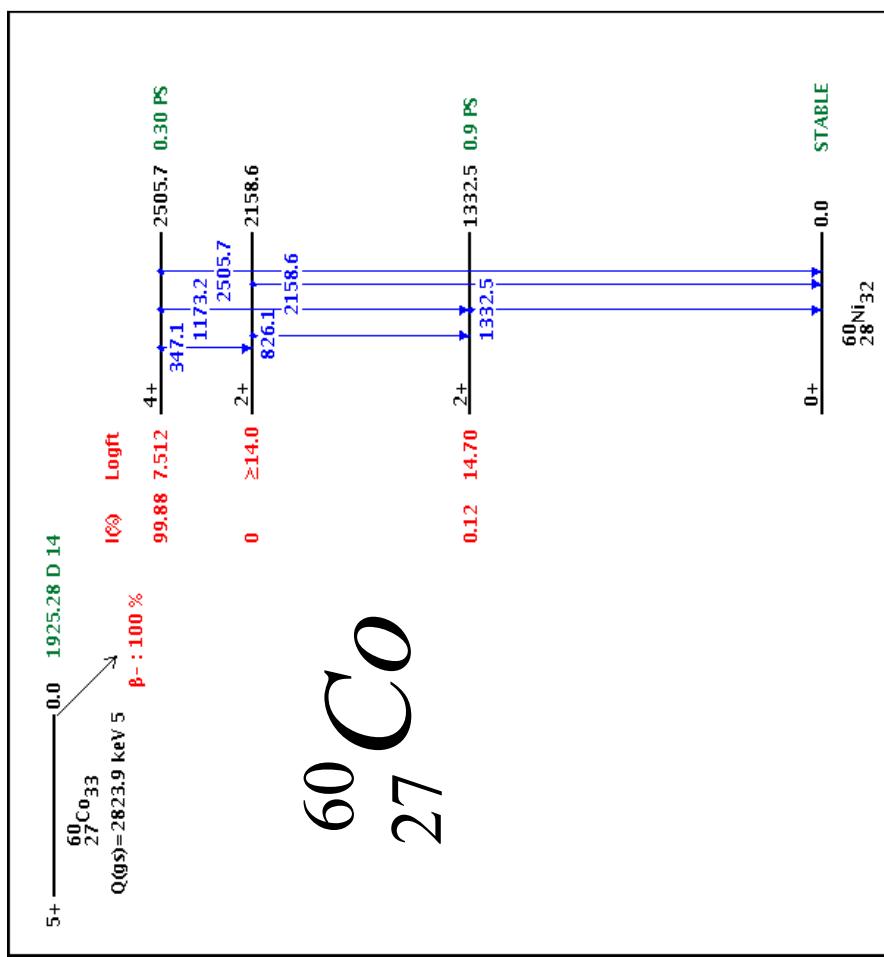
$$\left( cf : T_k = \frac{p^2}{2m} \right)$$

# Vzbuzené stavy jader a gamma záření





## Příklady známých gamu zářiců





## The Nobel Prize in Physics 1961

"for his pioneering studies of electron scattering in atomic nuclei and for his discoveries concerning the structure of the nucleons"

"for his researches concerning the resonance absorption of gamma radiation and his discovery in this connection of the effect which bears his name"

# Měření šířek gama přechodů a Mössbauerův jev

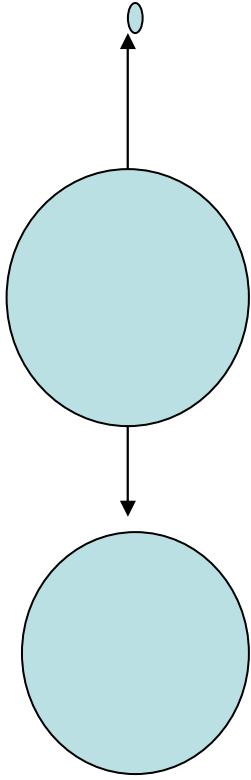
Při gama přechodu jádro odnáší část energie a tak je energie gama o něco menší



**Robert Hofstadter**  
1/2 of the prize  
USA

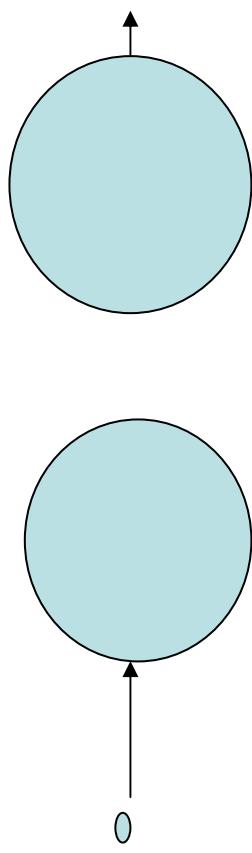


**Rudolf Ludwig Mössbauer**  
1/2 of the prize  
Federal Republic of Germany



$$\frac{E'_{\gamma}}{E_{\gamma}} = \frac{(M + E_{\gamma})^2 - M^2}{2(M + E_{\gamma})} = \frac{2ME_{\gamma} + E^2_{\gamma}}{2(M + E_{\gamma})} =$$
$$\frac{2E_{\gamma}(M + E_{\gamma}) - E^2_{\gamma}}{2(M + E_{\gamma})} = E_{\gamma} - \frac{E^2_{\gamma}}{2(M + E_{\gamma})}$$

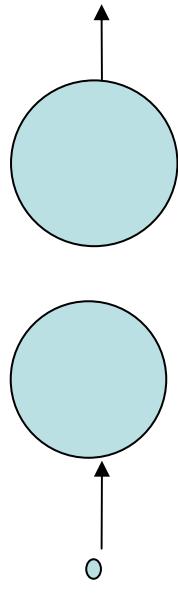
Naopak pro vybuzení jádra do excitovaného stavu je třeba energie o něco málo větší než odpovídá energii gamy přechodu



$$\begin{aligned}E''_\gamma + M &\rightarrow (M + E_\gamma) \\(E''_\gamma + M)^2 - (E''_\gamma)^2 &= (M + E_\gamma)^2 \\2E''_\gamma M &= 2E_\gamma M + E_\gamma^2 \\E''_\gamma &= E_\gamma + \frac{E_\gamma^2}{2M}\end{aligned}$$

$$(M + E_\gamma) \rightarrow M + E'_\gamma$$

$$E'_\gamma = E_\gamma - \frac{E^2_\gamma}{2(M + E_\gamma)}$$



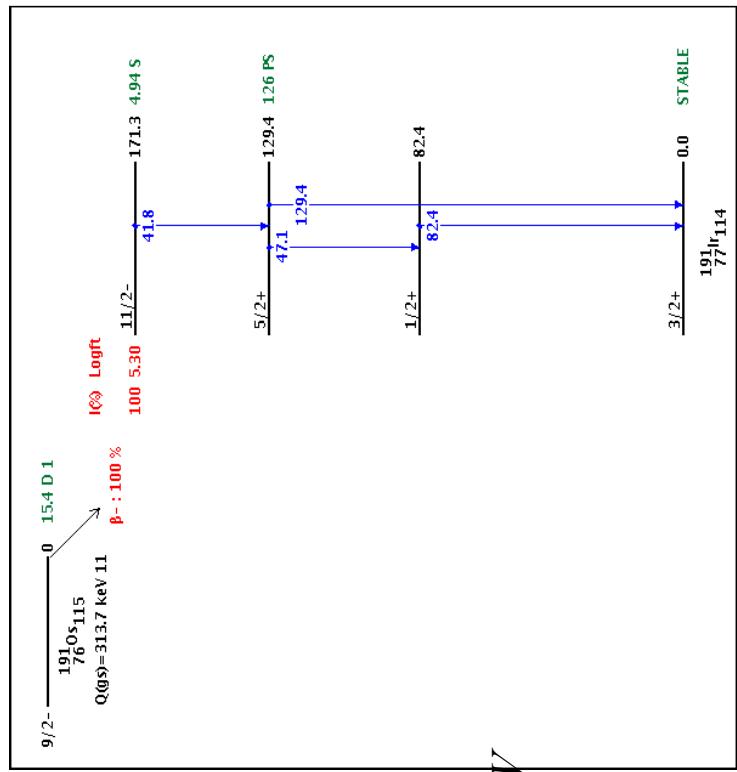
$$(^{191}_{77}Ir + 129keV) \rightarrow + E'_\gamma$$

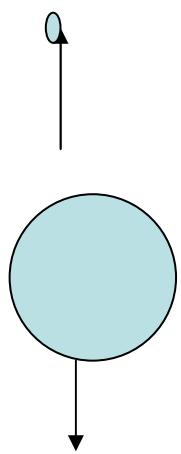
$$M_{^{191}_{77}Ir} = (77 \cdot 938,27 + (191 - 77) \cdot 939,57 - 191 \cdot 8,5) MeV =$$

$$177734 MeV$$

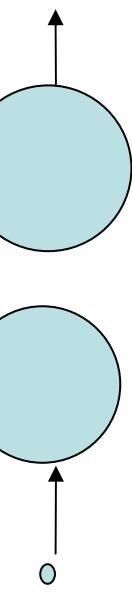
$$E''_\gamma = E_\gamma + M \rightarrow (M + E_\gamma)$$

$$E''_\gamma = E_\gamma + \frac{E^2_\gamma}{2M}$$



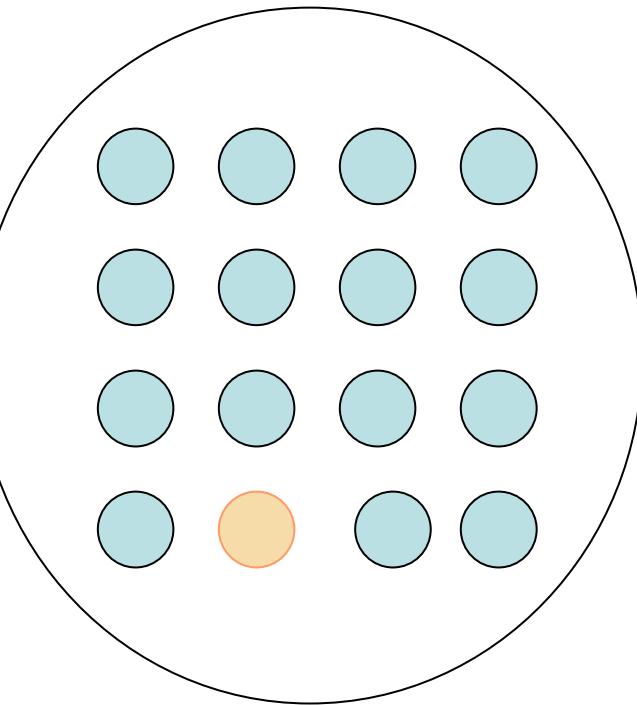
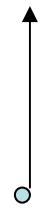
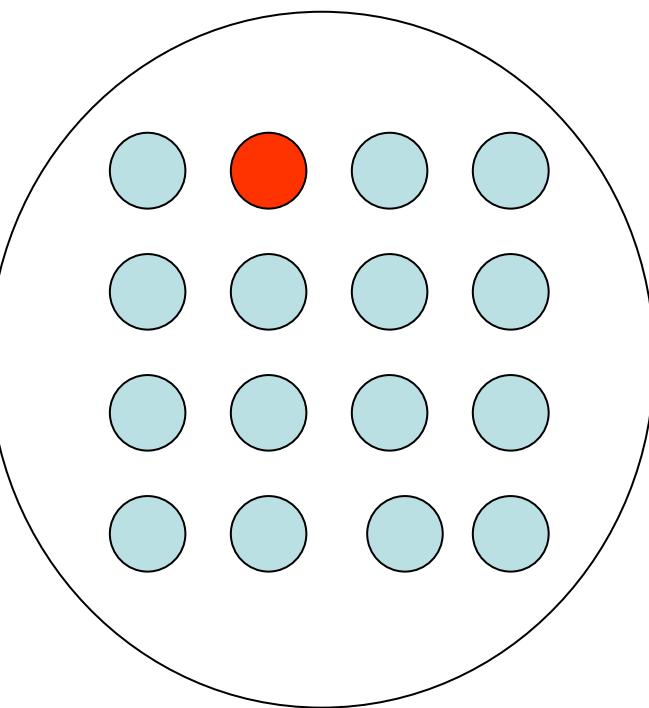


$$E' = E_\gamma - \frac{E^2 \gamma}{2(M + E_\gamma)}$$



$$E''_\gamma = E_\gamma + \frac{E^2 \gamma}{2M}$$

Je-li atom zabudován do krystalu ( $M \rightarrow M$  krystal), pak při nízké teplotě je odražen celý krystal a energie gama je správná na to aby v jiném takovém krystalu byla absorbována



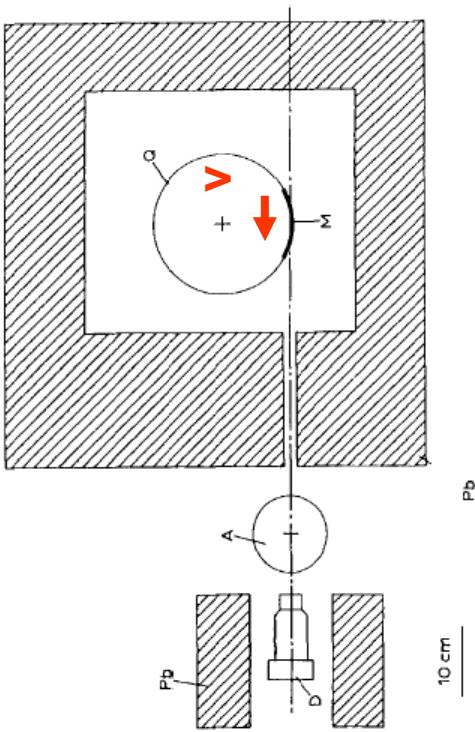


Fig. 7. Experimental arrangement. A, absorber-cryostat; Q, rotating cryostat with source; D, scintillation detector.

$$(^{191}\text{Ir} + 129\text{keV}) \rightarrow ^{191}\text{Ir}' + E_\gamma$$

$$\Gamma = 3,6 \cdot 10^{-6} \text{ eV}$$

$$129 \cdot 10^3 \beta \text{eV} = 3,6 \cdot 10^{-6} \text{ eV}$$

$$\beta = \frac{3,6 \cdot 10^{-6} \text{ eV}}{129 \cdot 10^3 \text{ eV}} = 2,8 \cdot 10^{-11}$$

$$\nu = 2,8 \cdot 10^{-11} \cdot 3 \cdot 10^{10} \text{ cm/s} = 0,84 \text{ cm/s}$$

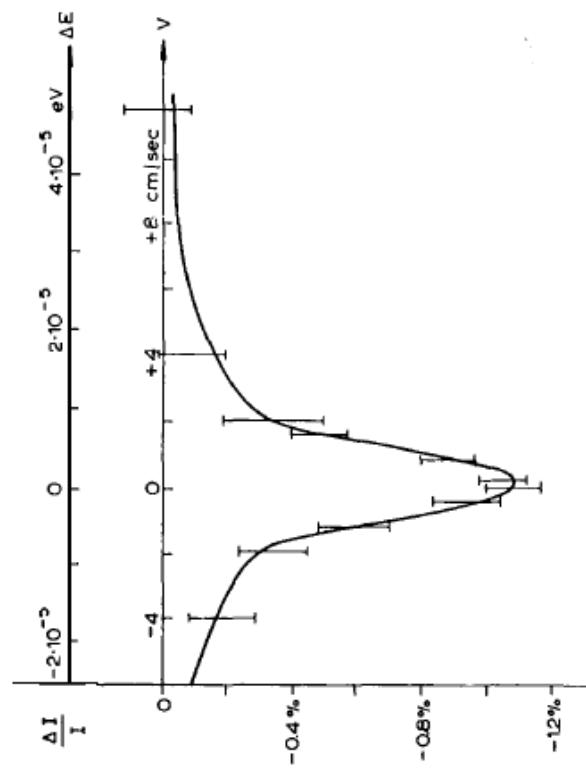
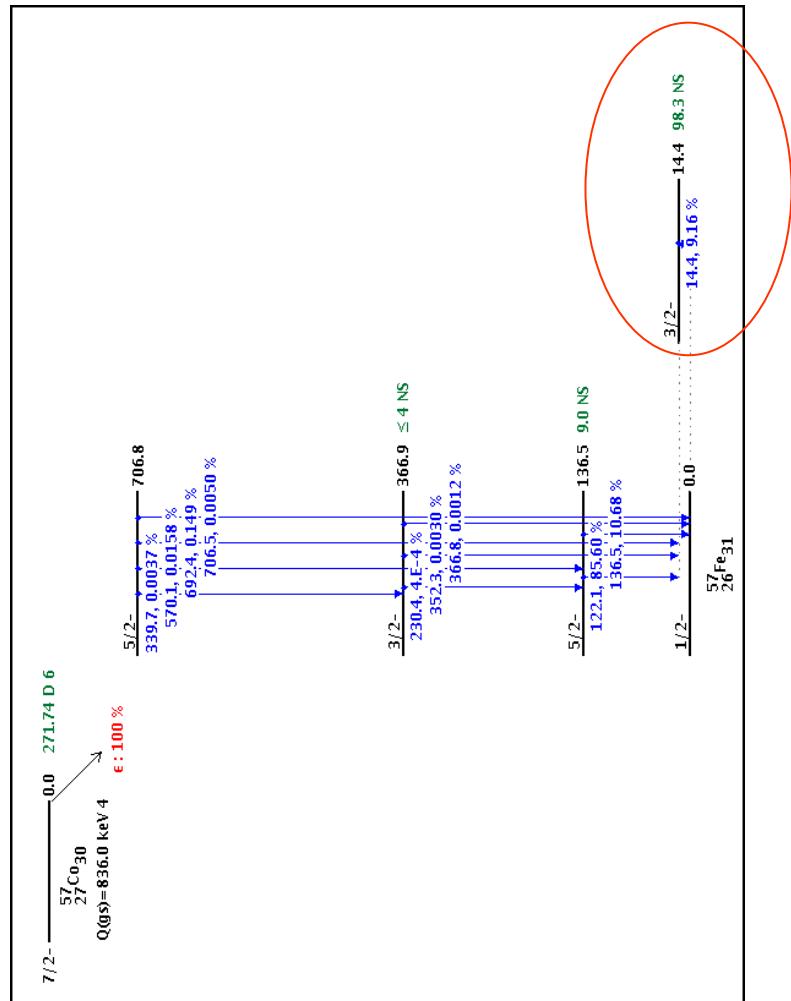


Fig. 8. Relative intensity ratio  $\Delta I/I$  of gamma radiation measured behind the resonant iridium absorber, in comparison with intensities measured behind a nonresonant absorber.

$^{57}_{26}Fe$

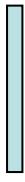
$$M_{^{57}_{26}Fe} = (26 \cdot 938,27 + (57 - 26) \cdot 939,57 - 57 \cdot 8,8) MeV = \\ 53020 MeV$$

$$\Gamma = \frac{\hbar c}{c \tau} = \frac{\hbar c}{c(T_{1/2} / \ln(2))} = \frac{197 \cdot 10^6 eV 10^{-13} cm}{30 cm / ns \cdot 98,3 ns / 0,69} = 46 \cdot 10^{-10} eV$$



**Detektor**

# Prověrka obecné teorie relativity



**Absorptor**



$\beta$

Rezonanční absorbce 14,4 keV gama. Rudý posuv gama je kompenzován pohybem absorbátoru. Experiment Repty

$$E'_{\gamma} = E_{\gamma} - \frac{E_{\gamma}}{c^2} gh$$

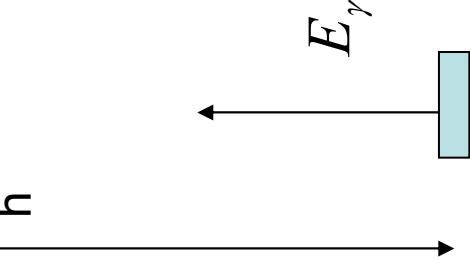
$$E''_{\gamma} = \gamma(E'_{\gamma} + \beta E'_{\gamma}) = \frac{1}{\sqrt{1-\beta^2}}(E'_{\gamma} + \beta E'_{\gamma}) \cong (1 + \frac{\beta^2}{2})(E'_{\gamma} + \beta E'_{\gamma}) \cong E'_{\gamma}(1 + \beta)$$

h

$$E_{\gamma}(1 - \frac{gh}{c^2})(1 + \beta) \frac{1}{\sqrt{1-\beta^2}} \cong E_{\gamma}(1 - \frac{gh}{c^2})(1 + \beta) \cong E_{\gamma}(1 - \frac{gh}{c^2} + \beta) = E_{\gamma}$$

$$\beta = \frac{gh}{c^2} = \frac{10m/s^2 \cdot 20m}{9 \cdot 10^{16} m^2/s^2} = 2,2 \cdot 10^{-15} \Rightarrow v = 2,2 \cdot 10^{-15} \cdot 3 \cdot 10^8 m/s =$$

$$6,6 \cdot 10^{-4} mm/s = 0,66 \mu m/s$$



## **Interakce jaderného záření s prostředím**

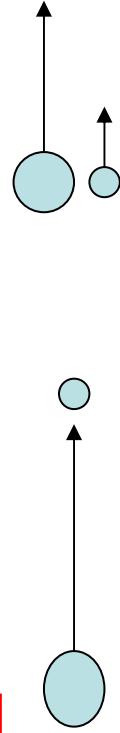
**Ionizace – nabité částice**  
**Brzdné záření – elektrony**  
**Interakce gamma záření**

## Ionizační ztráty

- týkají se všech nabitéch částic

**Tmax**

$$(E, \vec{P}) + (m_e, \vec{0}) = (E', \vec{P}') + (E_e, \vec{p}_e)$$



$$E + m_e = E' + E_e ; \quad \vec{P} = \vec{P}' + \vec{p}_e$$

$$E' = E + m_e - E_e ; \quad \left| \vec{P}' \right| = \left| \vec{P} \right| - \left| \vec{p}_e \right|$$

$$M^2 = E'^2 - \left| \vec{P}' \right|^2 = (E + m_e - E_e)^2 - \left( \left| \vec{P} \right| - \left| \vec{p}_e \right| \right)^2$$

$$M^2 = \left( E^2 - \left| \vec{P} \right|^2 \right) + \left( E_e^2 - \left| \vec{p}_e \right|^2 \right) + m_e^2 + 2Em_e - 2EE_e - 2m_eE_e + 2\left| \vec{P} \right|\left| \vec{p}_e \right|$$

$$0 = 2m_e^2 + 2Em_e - 2EE_e - 2m_eE_e + 2\left| \vec{P} \right|\left| \vec{p}_e \right|$$

$$\left( \left| \vec{P} \right| \left| \vec{p}_e \right| \right)^2 = (E_e - m_e)^2 (E + m_e)^2$$

$$\left| \vec{P} \right|^2 (E_e^2 - m_e^2) = \left| \vec{P} \right|^2 (E_e - m_e)(E + m_e) = (E_e - m_e)^2 (E + m_e)^2$$

$$\left| \vec{P} \right|^2 (\tau_{k,\max} + 2m_e) = (E + m_e)^2 \tau_{k,\max} ; \quad \tau_{k,\max} ((E + m_e)^2 - \left| \vec{P} \right|^2) = \left| \vec{P} \right|^2 2m_e$$

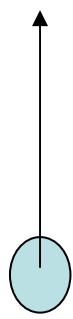
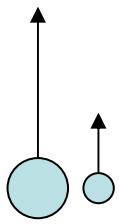
$$\tau_{k,\max} = \frac{2m_e \left| \vec{P} \right|^2}{M^2 + 2Em_e + m_e^2} = \frac{2m_e \left| \vec{P} \right|^2 / M^2}{1 + 2(E/M)(m_e/M) + (m_e/M)^2}$$

$$\tau_{k,\max} = \frac{2m_e \beta^2 \gamma^2}{1 + 2\gamma \frac{m_e}{M} + \left( \frac{m_e}{M} \right)^2}$$

**T<sub>max</sub>**

$$\tau_{k,\max} = \frac{2m_e\beta^2\gamma^2}{1 + 2\gamma \frac{m_e}{M} + \left(\frac{m_e}{M}\right)^2}$$

**T<sub>min</sub>**



## Ionizační ztráty - týkají se všech nabitéch částic

$$\frac{\pi(r^2_{\min}/8)}{\sin^4(\theta/2)} = \frac{d\sigma}{d\cos(\theta)} = \frac{d\sigma}{d(1-\cos(\theta))} = \frac{2P^2 d\sigma}{d(4P^2(1-\cos(\theta))/2)} = \frac{2P^2 d\sigma}{d(4P^2 \sin^2(\theta/2))} = \frac{2P^2 d\sigma}{dq^2}$$

$$\frac{2P^2 d\sigma}{dq^2} = \frac{\pi(r^2_{\min}/8)}{\sin^4(\theta/2)} = \frac{(2P)^4 \pi(r^2_{\min}/8)}{(2P)^4 \sin^4(\theta/2)} = \frac{(2P)^4 \pi(r^2_{\min}/8)}{q^4}$$

$$\frac{d\sigma}{dq^2} = \frac{P^2 \pi r^2_{\min}}{q^4} \Rightarrow \frac{d\sigma}{dt} = \frac{P^2 \pi r^2_{\min}}{|t|^2}$$

$$r_e = \frac{\alpha(\hbar c)}{m_e}$$

$$\frac{d\sigma}{2m_e d\frac{|t|}{2m_e}} = \frac{P^2 \pi r^2_{\min}}{4m_e^2 \left(\frac{|t|}{2m_e}\right)^2}$$

$$\frac{d\sigma}{d\tau_k} = \frac{P^2 \pi r^2_{\min}}{2m_e \tau^2_k} = \frac{m_e P^2 \pi \alpha^2 (\hbar c)^2}{2m_e^2 T_k^2} \frac{1}{\tau^2_k} = \frac{P^2}{2T_k^2} m_e \pi r^2 e \frac{1}{\tau^2_k} =$$

$$\frac{(\gamma\beta)^2 M^2}{2(\gamma-1)^2} \frac{1}{m_e \pi r^2 e \tau^2_k} = \frac{\gamma^2 - 1}{2(\gamma-1)^2} \frac{1}{m_e \pi r^2 e \tau^2_k} = \frac{\gamma+1}{2(\gamma-1)} \frac{1}{m_e \pi r^2 e} \frac{1}{\tau^2_k} =$$

$$\frac{(\gamma+1)^2}{2(\gamma^2-1)} \frac{1}{m_e \pi r^2 e \tau^2_k} = \frac{(\gamma+1)^2}{2(\gamma\beta)^2} \frac{1}{m_e \pi r^2 e \tau^2_k} = 2 \left( \frac{\gamma+1}{2} \right)^2 \frac{1}{(\gamma\beta)^2} \frac{1}{m_e \pi r^2 e} \frac{1}{\tau^2_k}$$

$$\frac{d\sigma}{d\tau_k} = 2\left(\frac{\gamma+1}{2}\right)^2 \frac{1}{(\gamma\beta)^2} m_e \pi r_e^2 \frac{1}{\tau_k^2}$$

$$p=\sigma\cdot\Delta x\cdot\rho\cdot\frac{N_A}{A}Z$$

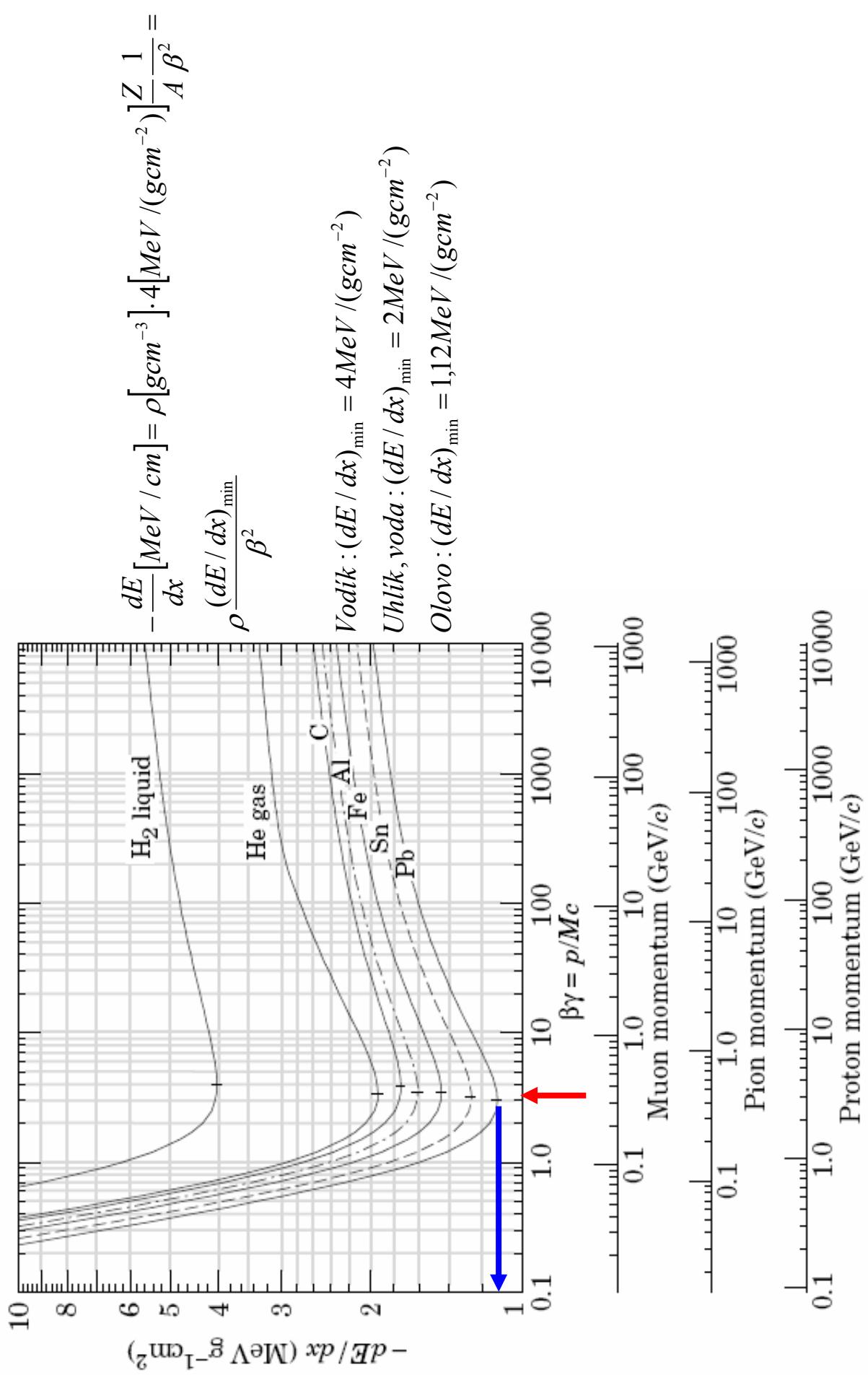
$$\frac{dp}{d\tau_k} = \Delta x \cdot \rho \cdot \frac{N_A}{A} Z \cdot 2\left(\frac{\gamma+1}{2}\right)^2 \frac{1}{(\gamma\beta)^2} m_e \pi r_e^2 \frac{1}{\tau_k^2}$$

$$\langle \tau_k \rangle = \int_{\tau_{\min}}^{\tau_{\max}} \tau_k \frac{dp}{d\tau_k} d\tau_k = \Delta x \cdot \rho \cdot \frac{N_A}{A} Z \cdot 2\left(\frac{\gamma+1}{2}\right)^2 \frac{1}{(\gamma\beta)^2} m_e \pi r_e^2 \int_{\tau_{\min}}^{\tau_{\max}} \tau_k \frac{1}{\tau_k^2} d\tau_k =$$

$$\Delta x \cdot \rho \cdot \frac{N_A}{A} Z \cdot 2\left(\frac{\gamma+1}{2}\right)^2 \frac{1}{(\gamma\beta)^2} m_e \pi r_e^2 \ln\left(\frac{\tau_{\max}}{\tau_{\min}}\right)$$

$$-\frac{\Delta T_k}{\Delta x} = \rho \cdot \frac{N_A}{A} Z \cdot 2\left(\frac{\gamma+1}{2}\right)^2 \frac{1}{(\gamma\beta)^2} m_e \pi r_e^2 \ln\left(\frac{\tau_{\max}}{\tau_{\min}}\right)$$

$$-\frac{dE}{dx} = K_z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$



## Doběh:

$$-\frac{dE}{dx} = \frac{(dE/dx)_{\min} \cdot \rho}{\beta^2} = \frac{(dE/dx)_{\min} \cdot \rho}{(E^2 - m^2)/E^2}$$

$$-dE \frac{E^2 - m^2}{E^2} = (dE/dx)_{\min} \cdot \rho \cdot dx$$

$$-\int_{E_0}^m dE \left(1 - \frac{m^2}{E^2}\right) = (dE/dx)_{\min} \cdot \rho \int_0^R dx$$

$$-(m - E_0) - m^2 \left(\frac{1}{m} - \frac{1}{E_0}\right) = (dE/dx)_{\min} \cdot \rho \cdot R$$

$$E_0 - 2m + \frac{m^2}{E_0} = (dE/dx)_{\min} \cdot \rho \cdot R$$

$$R = \boxed{\frac{E_0 - m}{(dE/dx)_{\min} \cdot \rho} \frac{E_0 - m}{E_0}} = \boxed{\frac{T_k}{(dE/dx)_{\min} \cdot \rho} \frac{T_k}{T_k + m}}$$

$T_k << m$ :

$$R = \frac{T_k}{(dE/dx)_{\min} / \rho} \frac{T_k}{m} \approx T_k^2$$

$T_k >> m$ :

$$R = \frac{T_k}{(dE/dx)_{\min} / \rho} \approx T_k$$

*Př.:*

$$\alpha, m = (2 \cdot 938,27 + 2 \cdot 939,57 - 28) MeV = 3725,68 MeV$$

$$T_k = 7,7 MeV$$

*vzduch:*

$$(dE / dx)_{\min} = 2,0 MeV / (g \cdot cm^{-2})$$

$$\rho = 0,0012 g \cdot cm^{-3}$$

$$R = \frac{T_k}{(dE / dx)_{\min} \cdot \rho} = \frac{T_k}{2,0 \frac{MeV}{g \cdot cm^{-2}} \cdot 0,0012 g \cdot cm^{-3}} = \frac{7,7 MeV}{(7,7 + 3725,68) MeV} = 6,62 cm$$

*voda:*

$$(dE / dx)_{\min} = 2,0 MeV / (g \cdot cm^{-2})$$

$$\rho = 1,0 g \cdot cm^{-3}$$

$$R = \frac{T_k}{(dE / dx)_{\min} \cdot \rho} = \frac{T_k}{2,0 \frac{MeV}{g \cdot cm^{-2}} \cdot 1,0 g \cdot cm^{-3}} = \frac{7,7 MeV}{(7,7 + 3725,68) MeV} = 7,9 \cdot 10^{-3} cm = 79 \mu m$$

*zlato:*

$$(dE / dx)_{\min} = 1,13 MeV / (g \cdot cm^{-2})$$

$$\rho = 20 g \cdot cm^{-3}$$

$$R = \frac{T_k}{(dE / dx)_{\min} \cdot \rho} = \frac{T_k}{1,13 \frac{MeV}{g \cdot cm^{-2}} \cdot 20,0 g \cdot cm^{-3}} = \frac{7,7 MeV}{(7,7 + 3725,68) MeV} = 7,0 \cdot 10^{-4} cm = 7 \mu m$$

Interakce elektronů:

Jakož i jiné nabité částice ionizují,

$$T_{k,\max} = \frac{2m_e\beta^2\gamma^2}{1 + 2\gamma \frac{m_e}{M} + \left(\frac{m_e}{M}\right)^2}$$

$$M = m_e$$

$$T_{k,\max} = \frac{2m_e\beta^2\gamma^2}{1 + 2\gamma + 1} = \frac{2m_e\beta^2\gamma^2}{2(1 + \gamma)} = \frac{m_e(\gamma^2 - 1)}{(1 + \gamma)} = m_e(\gamma - 1) = T_k$$

Při energiích  $E > 20 \text{ MeV}$  dominuje brzdné záření – bremssstrahlung

$$-\frac{dE}{dx} = \frac{E}{X_0}$$

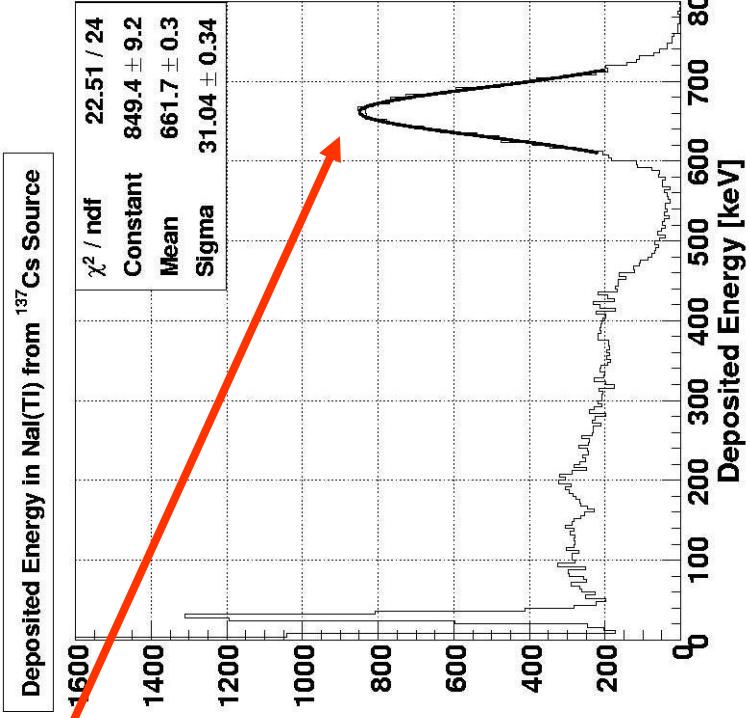
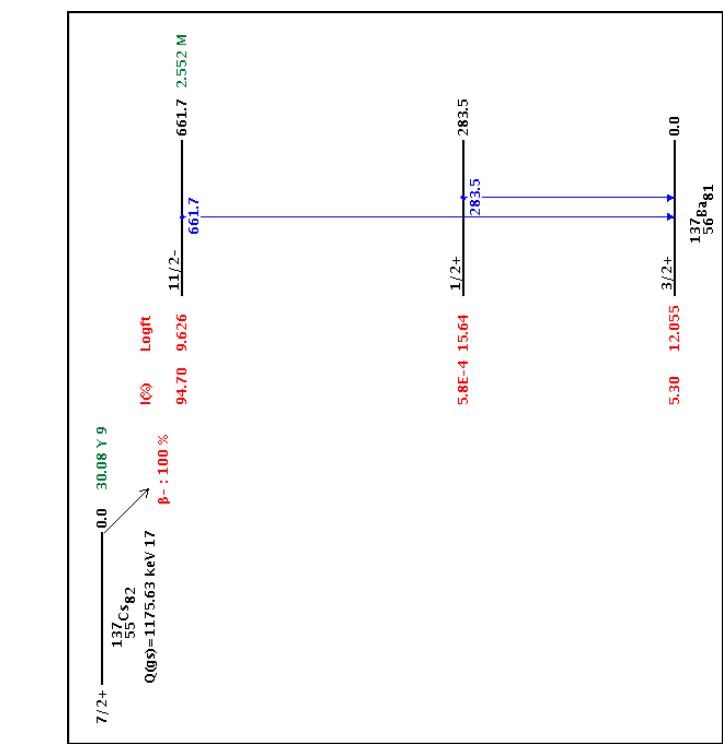
$$E = E_0 e^{-L/X_0}$$

$$X_0(\text{vzduch}) = 304 \text{ m}; X_0(\text{voda}) = 36 \text{ cm}$$

$$X_0(\text{Fe}) = 1,8 \text{ cm}; X_0(\text{Pb}) = 5,6 \text{ mm}$$

# Interakce záření gamma

Fotoefekt, Comptonův rozptyl, Tvorba páru elektron-positron



Compton :

$$E_\gamma + m_e = E'_\gamma + \sqrt{m_e^2 + p_e^2} \Rightarrow p_e^2 = (E_\gamma - E'_\gamma + m_e)^2 - m_e^2$$

$$E_\gamma = E'_\gamma \cos(\theta) + p_e \cos(\phi)$$

$$E'_\gamma \sin(\theta) = p_e \sin(\phi)$$

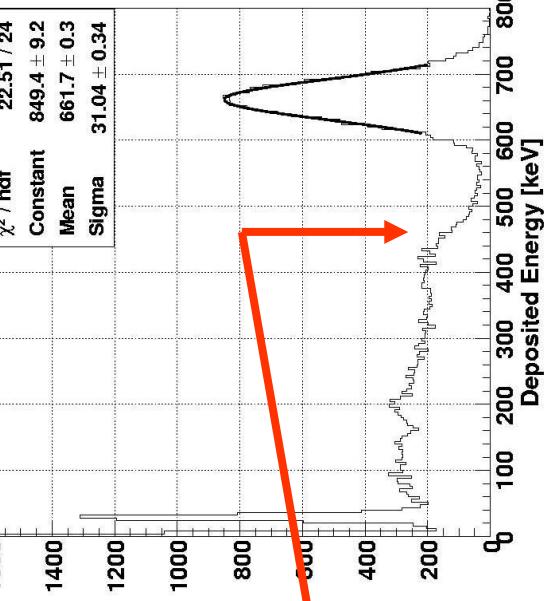
$$\Rightarrow p_e^2 = (E_\gamma - E'_\gamma \cos(\theta_\gamma))^2 + (E'_\gamma \sin(\theta_\gamma))^2$$

$$(E_\gamma - E'_\gamma \cos(\theta))^2 + (E'_\gamma \sin(\theta))^2 = (E_\gamma - E'_\gamma + m_e)^2 - m_e^2$$

$$E'^2_\gamma - 2E'_\gamma E'_\gamma \cos(\theta) + E'^2_\gamma = E'^2_\gamma - 2E'_\gamma E'_\gamma + 2m_e E'_\gamma - 2E'_\gamma m_e$$

$$E'_\gamma (E_\gamma (1 - \cos(\theta)) + m_e) = m_e E'_\gamma$$

$$E'_\gamma = \frac{E'_\gamma}{1 + \frac{E'_\gamma}{m_e} (1 - \cos(\theta))}$$



$$\Delta E_\gamma = E_\gamma - E'_\gamma = E_\gamma - \frac{E_\gamma}{1 + \frac{E_\gamma}{m_e} (1 - \cos(\theta))} = \frac{E_\gamma (1 - \cos(\theta))}{1 + \frac{E_\gamma}{m_e} (1 - \cos(\theta))} E_\gamma$$

$$\Rightarrow \Delta E_{\gamma_{\max}} = E_\gamma \frac{2 \frac{E_\gamma}{m_e}}{1 + 2 \frac{E_\gamma}{m_e}} = \frac{2 E_\gamma^2}{m_e + 2 E_\gamma}$$

# Tvorba elektron-positronových páru

$$\gamma + M \rightarrow M + e^+ + e^-$$

Prahová energie gama:

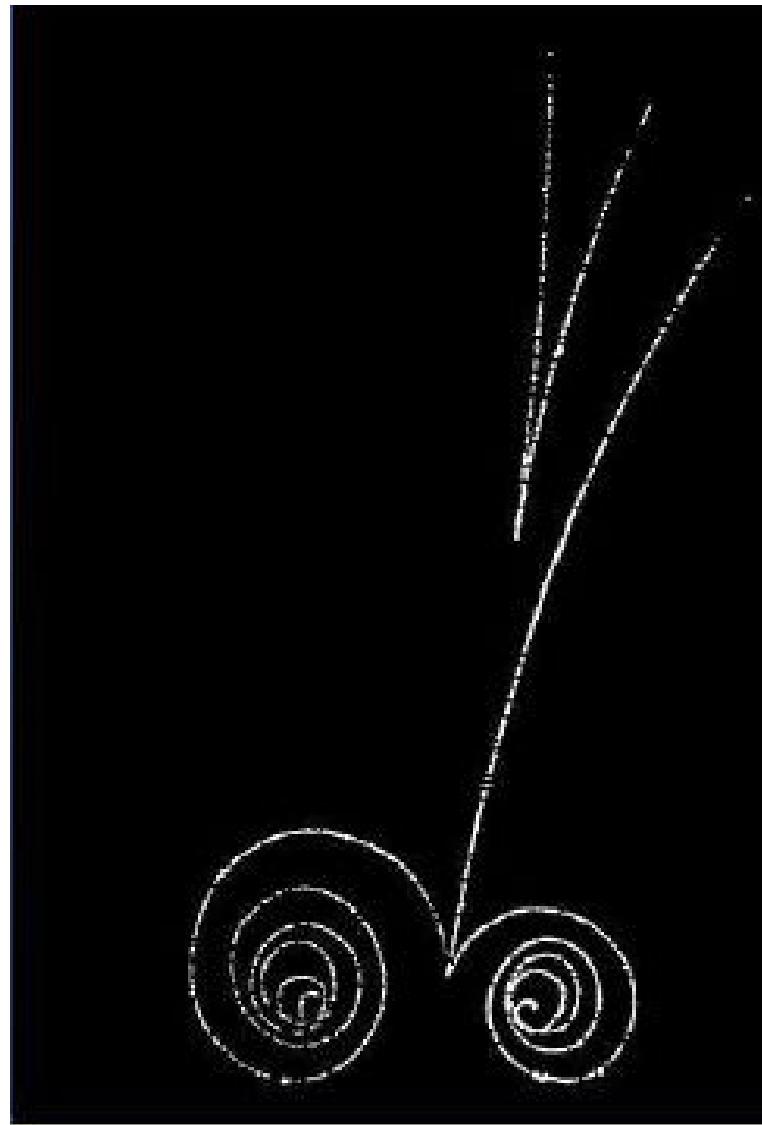
$$(E_{\gamma}^{thr} + M)^2 - E_{\gamma}^{thr2} = (M + 2m_e)^2$$

$$2E_{\gamma}^{thr}M = 4m_eM + 4m_e^2$$

$$E_{\gamma}^{thr} = 2m_e \left( 1 + \frac{m_e}{M} \right)$$

$$M \gg m_e \Rightarrow E_{\gamma}^{thr} \approx 2m_e$$

$$M = m_e \Rightarrow E_{\gamma}^{thr} = 4m_e$$



$$P(x) = 1 - e^{-(\sigma_{fotoeffekt} + \sigma_{Compton} + \sigma_{pairs})x \frac{N_A}{A} \rho}$$

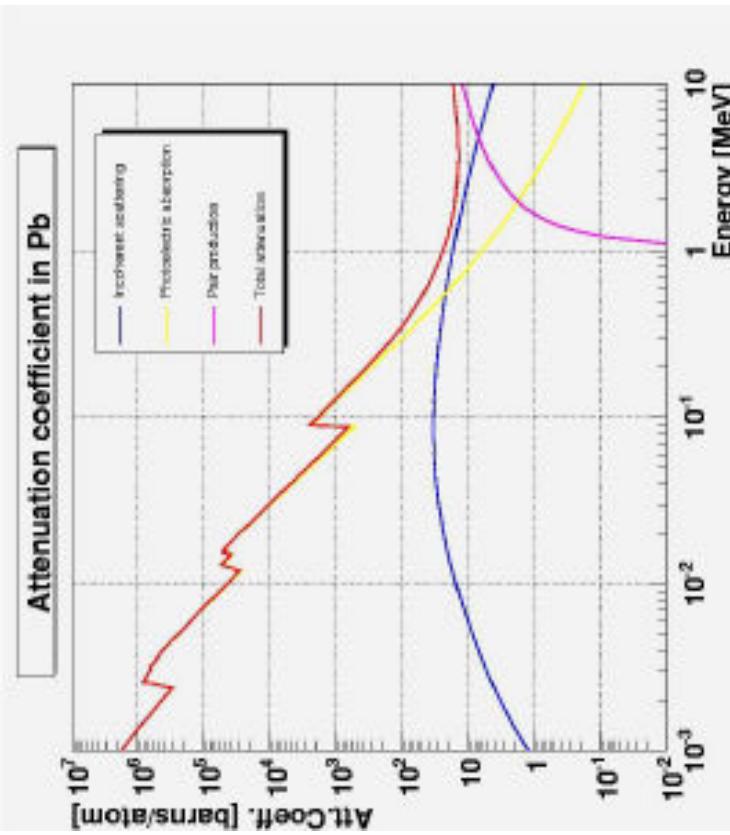
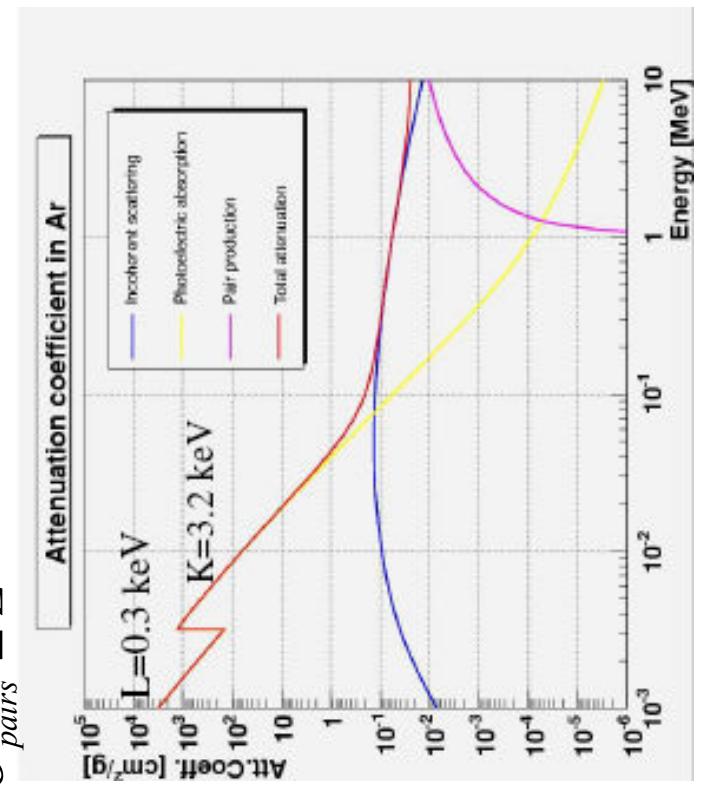
$$I(x) = I_0(1 - P(x)) = I_0 e^{-(\sigma_{fotoeffekt} + \sigma_{Compton} + \sigma_{pairs})x \frac{N_A}{A}} = I_0 e^{-\mu \alpha x}$$

$$\mu = \left( \sigma_{fotoeffekt} + \sigma_{Compton} + \sigma_{pairs} \right) \frac{N_A}{A}$$

$$\sigma_{fotoeffekt} \approx Z^{(4-5)}$$

$$\sigma_{Compton} \approx Z$$

$$\sigma_{pairs} \approx Z^2$$

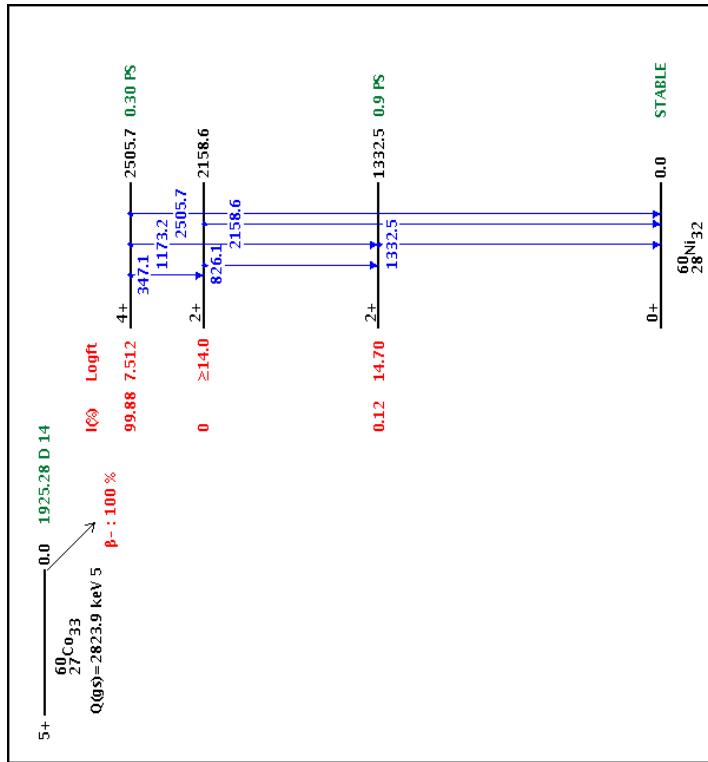


# Radiace, jednotky

Aktivita se měří v Becquerel Bq = 1 rozpad/s

Dávka Gray G = 1 J/kg

Sievert Sv= Gray x Quality factor – určuje jak je dané záření nebezpečné pro člověka. Nejméně nebezpečné jsou miony, gama, beta (Q=1), nejvíce alfa a ostatní těžké ionty (Q=20) a neutrony s malou energií (Q~5).  
Povoleno je 50mSv za rok a 100 mSv za pět let.



1MBq kobaltového gama záříče:

$$-\Delta N = N \frac{\Delta t}{T_{1/2} / \ln(2)}$$

$$-\Delta N / \Delta t = A = \frac{N}{T_{1/2} / \ln(2)}$$

$$N = A[Bq = \text{s}^{-1}] \cdot T_{1/2} [\text{s}] / \ln(2)$$

$$A = 1MBq; T_{1/2} = 1925 days$$

$$N = 10^6 \text{ s}^{-1} 1925 \cdot 86400 \text{ s} / 0,69 = 0,24 \cdot 10^{15} \text{ } ^{60}_{27}\text{Co}$$

$$60gCo = 6,023 \cdot 10^{23} atoms$$

$$0,24 \cdot 10^{15} \text{ } ^{60}_{27}\text{Co} \Rightarrow 0,24 \cdot 10^{15} \cdot 60g / 6,023 \cdot 10^{23} = 2,4 \cdot 10^{-8} g = 24ng!!$$

## Dávky záření:

### Miony z kosmického záření:

$$(dE / dx)_{\min} = 2 \text{MeV} / g \cdot cm^{-2} \Rightarrow \rho = 1g / cm^3; (dE / dx)_{\min} = 2 \text{MeV} / cm$$

50cm

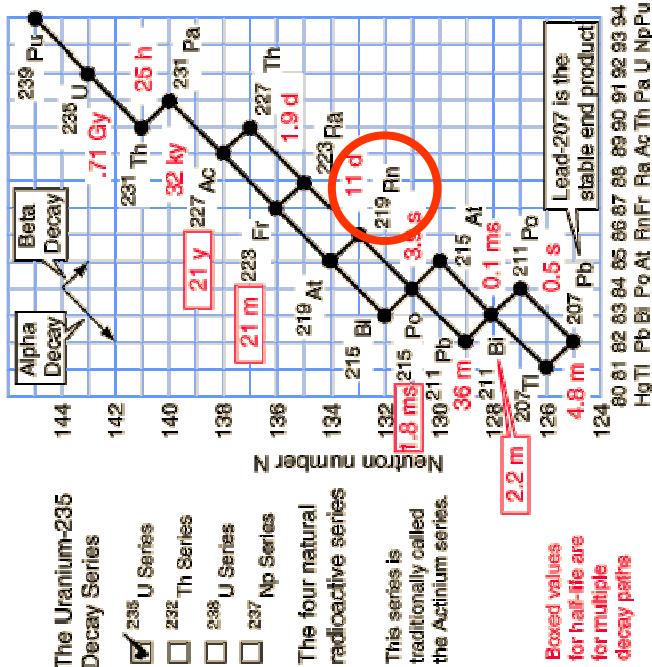
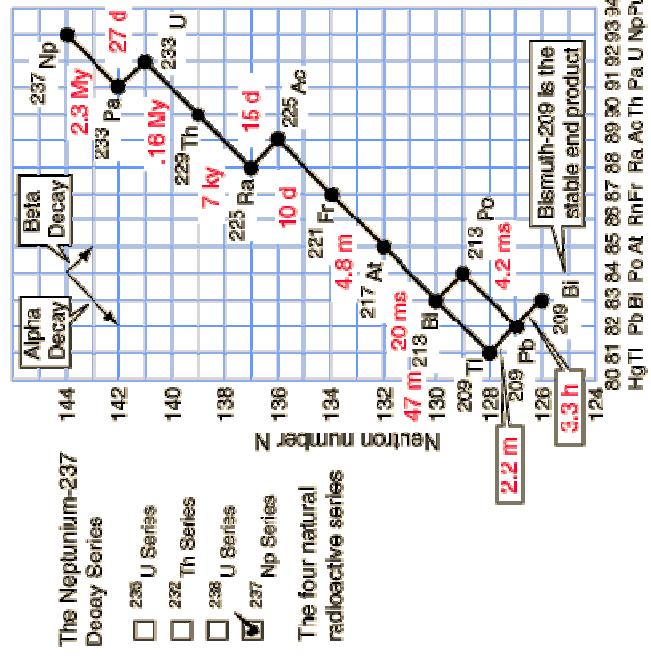
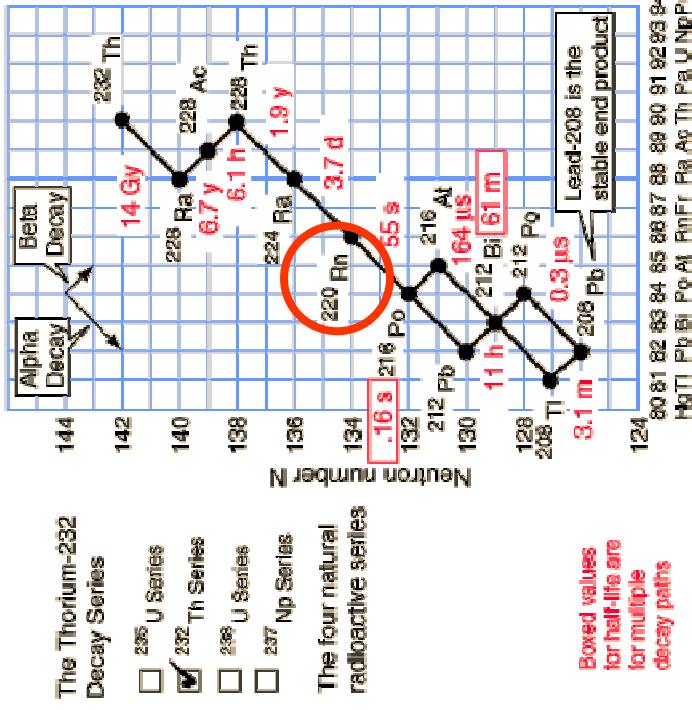
$$1mion / s / dm^2 \Rightarrow 6mion / s$$

$$DE = 2 \text{MeV} / cm \cdot 50cm \cdot 6mion / s \cdot 31,4 \cdot 10^7 s / year = 1,88 \cdot 10^{11} \text{MeV} / year =$$
$$1,88 \cdot 10^{11} \cdot 10^6 \cdot 1,6 \cdot 10^{-19} J / year = 3 \cdot 10^{-2} J / year$$
$$D = 3 \cdot 10^{-2} J / year / 80kg = 0,375mJ / year = \underline{\underline{0,375mSv / rok}}$$

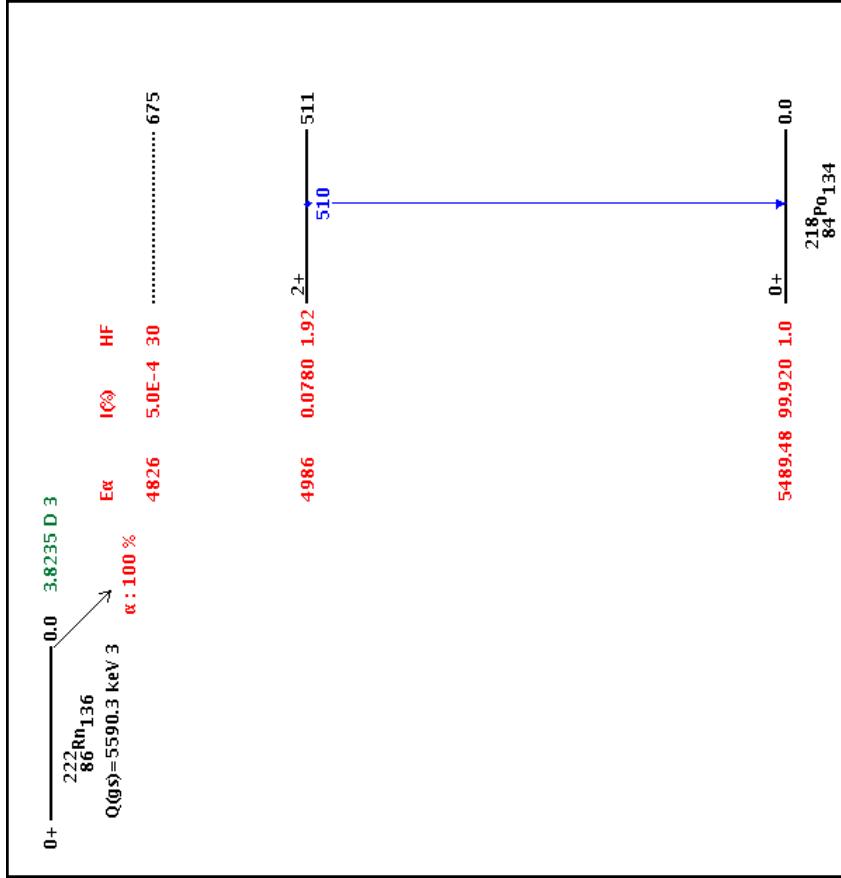
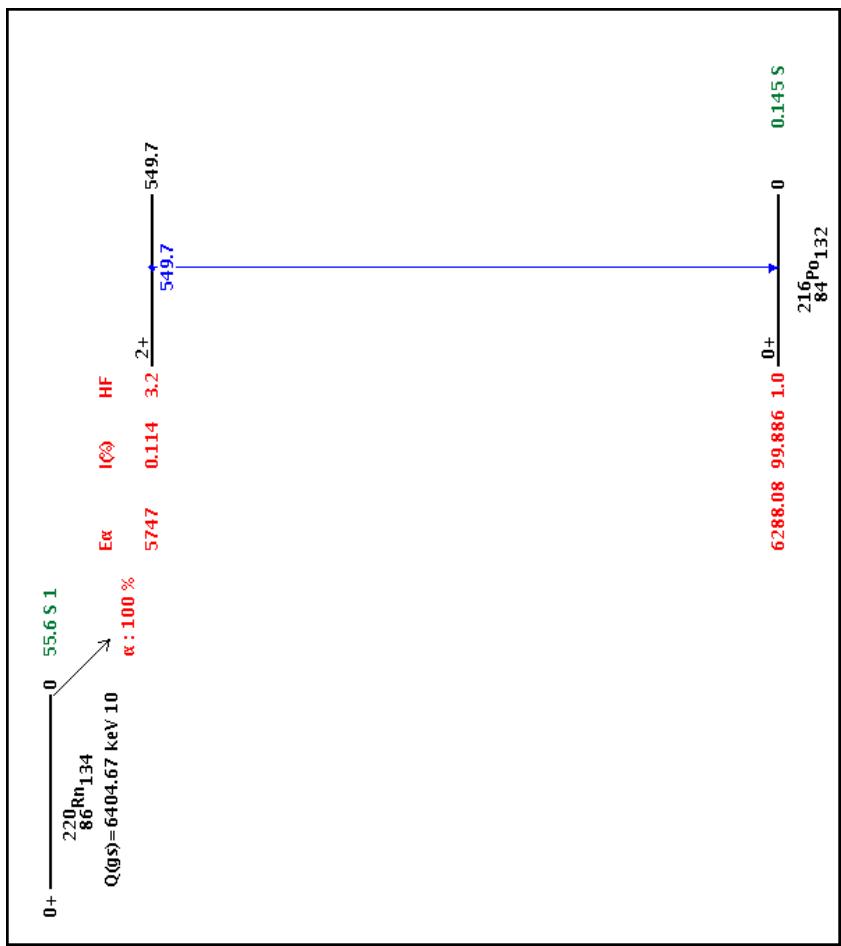
### V Troji ve vstupní hale:

$$D = 0,09 \mu Sv / hod \Rightarrow 0,09 \cdot 365 \cdot 24 \mu Sv / rok = 0,788mSv / rok$$

**Radon 2mSv/rok**



Radon je plyn, vdechujeme ho a alfa částice mohou poškodit plíce



# Následky ozáření:

## Table of exposure levels and symptoms

Dose-equivalents are presently stated in sieverts:

### **0.05–0.2 Sv (5–20 REM)**

No symptoms.

### **0.2–0.5 Sv (20–50 REM)**

No noticeable symptoms. Red blood cell count decreases temporarily.

### **0.5–1 Sv (50–100 REM)**

Mild radiation sickness with headache and increased risk of infection due to disruption of immunity cells. Temporary male sterility is possible.

### **1–2 Sv (100–200 REM)**

*Light radiation poisoning, 10% fatality after 30 days (LD 10/30).* Typical symptoms include mild to moderate nausea (50% probability at 2 Sv), with occasional vomiting, beginning 3 to 6 hours after irradiation and lasting for up to one day. This is followed by a 10 to 14 day latent phase, after which light symptoms like general illness and fatigue appear (50% probability at 2 Sv). The immune system is depressed, with convalescence extended and increased risk of infection. Permanent female sterility is common. Spontaneous abortion or stillbirth will occur in pregnant women.

### **2–3 Sv (200–300 REM)**

*Moderate radiation poisoning, 35% fatality after 30 days (LD 35/30).* Nausea is common (100% at 3 Sv), with 50% risk of vomiting at 2.8 Sv. Symptoms onset at 1 to 6 hours after irradiation and last for 1 to 2 days. After that, there is a 7 to 14 day latent phase, after which the following symptoms appear: loss of hair all over the body (50% probability at 3 Sv), fatigue and general illness. There is a massive loss of leukocytes (white blood cells), greatly increasing the risk of infection. Permanent female sterility is possible. Convalescence takes one to several months.

### **3–4 Sv (300–400 REM)**

*Severe radiation poisoning, 50% fatality after 30 days (LD 50/30).* Other symptoms are similar to the 2–3 Sv dose, with uncontrollable bleeding in the mouth, under the skin and in the kidneys (50% probability at 4 Sv) after the latent phase. 180 Sv (18,000 REM) to his upper body in an accident at Los Alamos, New Mexico, USA on 30 December 1958 survived for 36 hours.

# Následky ozáření:

## 4–6 Sv (400–600 REM)

*Acute radiation poisoning, 60% fatality after 30 days (LD 60/30).* Fatality increases from 60% at 4.5 Sv to 90% at 6 Sv (unless there is intense medical care). Symptoms start half an hour to two hours after irradiation and last for up to 2 days. After that, there is a 7 to 14 day latent phase, after which generally the same symptoms appear as with 3-4 Sv irradiation, with increased intensity. Female sterility is common at this point. Convalescence takes several months to a year. The primary causes of death (in general 2 to 12 weeks after irradiation) are infections and internal bleeding.

## 6–10 Sv (600–1,000 REM)

*Acute radiation poisoning, near 100% fatality after 14 days (LD 100/14).* Survival depends on intense medical care. Bone marrow is nearly or completely destroyed, so a bone marrow transplant is required. Gastric and intestinal tissue are severely damaged. Symptoms start 15 to 30 minutes after irradiation and last for up to 2 days. Subsequently, there is a 5 to 10 day latent phase, after which the person dies of infection or internal bleeding. Recovery would take several years and probably would never be complete.

Devair Alves Ferreira received a dose of approximately 7.0 Sv (700 REM) during the Goiânia accident and survived, partially due to his fractionated exposure.

## 10–50 Sv (1,000–5,000 REM)

*Acute radiation poisoning, 100% fatality after 7 days (LD 100/7).* An exposure this high leads to spontaneous symptoms after 5 to 30 minutes. After powerful fatigue and immediate nausea caused by direct activation of chemical receptors in the brain by the irradiation, there is a period of several days of comparative well-being, called the latent (or "walking ghost") phase [citation needed]. After that, cell death in the gastric and intestinal tissue, causing massive diarrhea, intestinal bleeding and loss of water, leads to water-electrolyte imbalance. Death sets in with delirium and coma due to breakdown of circulation. Death is currently inevitable; the only treatment that can be offered is pain therapy.

Louis Slotin was exposed to approximately 21 Sv in a criticality accident on 21 May 1946, and died nine days later on 30 May. At this dose the skin can be damaged. Here is a photo of a man who received a 10 to 20 Gy gamma whole body dose as a result of an industrial accident. He died about 10 days after the photo was taken, about 30 days after the event.

## More than 50 Sv (>5,000 REM)

A worker receiving 100 Sv (10,000 REM) in an accident at Wood River, Rhode Island, USA on 24 July 1964 survived for 49 hours after exposure, and an operator receiving between 60 and 180 Sv (18,000 REM) to his upper body in an accident at Los Alamos, New Mexico, USA on 30 December 1958 survived for 36 hours.