

Fyzika V

Rupert Leitner

Rupert.Leitner@mff.cuni.cz

ÚČJF MFF UK 838A, 221912444

Doporučená literatura:

W.S.C. Williams: Nuclear and Particle Physics

Přednáška 1 (2.10.2007)

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- 1.3. Atom vodíku – opakování
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1.1. Jednotky

$$E^2 = (mc^2)^2 + (pc)^2$$

E eV, MeV, GeV

$$mc^2 \longrightarrow m$$

$$pc \longrightarrow p$$

$$E^2 = m^2 + p^2$$

neutron

$$m_n = 939,57 \text{ MeV}$$

proton

$$m_p = 938,27 \text{ MeV}$$

electron

$$m_e = 0,511 \text{ MeV}$$

$$m_n - m_p = 1,3 \text{ MeV}$$

Coulombův zákon

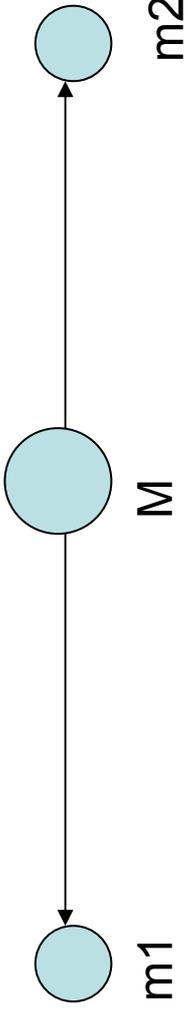
$$V(r) = \frac{e^2}{4\pi\epsilon_0 r} = \frac{e^2}{4\pi\epsilon_0} \frac{\hbar c}{\hbar c r} = \alpha \frac{\hbar c}{r}$$

α - konstanta jemné struktury

$$\alpha = \frac{1}{137,036} \approx \frac{1}{137}$$

$$\hbar c = 197,327 \text{ MeV} \cdot \text{fm}$$

1.2. Kinematika



$$p_1 = p_2 = p^*$$

$$\sqrt{(p^*)^2 + m_1^2} + \sqrt{(p^*)^2 + m_2^2} = M$$

$$p^* = \frac{\sqrt{M^2 - (m_1 + m_2)^2} \sqrt{M^2 - (m_1 - m_2)^2}}{2M}$$

$$m_1 = m_2 = 0; p^* = M / 2$$

$$m_1 = m_2 = m; p^* = \frac{\sqrt{M^2 - 4m^2}}{2} = \sqrt{M^2 / 4 - m^2}$$

$$m_1 = m, m_2 = 0; p^* = \frac{M^2 - m^2}{2M}$$

$$\pi^0 \rightarrow \gamma\gamma; m_1 = m_2 = 0; p^* = M_{\pi^0} / 2 = 135 \text{ MeV} / 2 = 67,5 \text{ MeV}$$

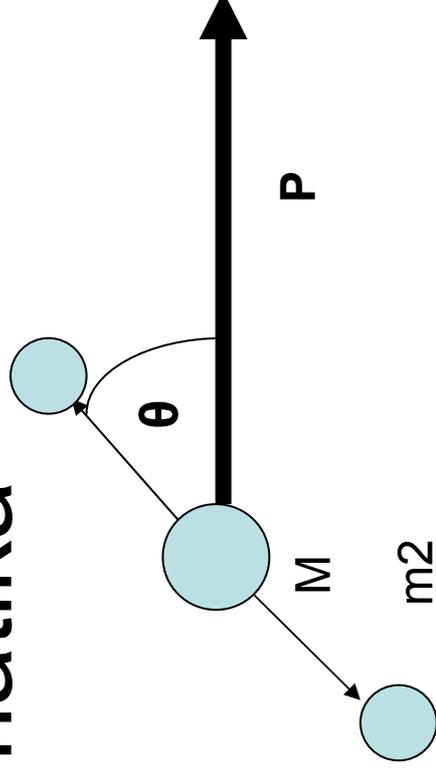
$$\rho^0 \rightarrow \pi^+ \pi^-; m_1 = m_2 = m_\pi; p^* = \sqrt{M_\rho^2 / 4 - m_\pi^2} = \sqrt{770^2 / 4 - 140^2} \approx 359 \text{ MeV}$$

$$\pi^+ \rightarrow \mu^+ \nu_\mu; m_1 = m_\mu, m_2 = 0; p^* = \frac{M_\pi^2 - m_\mu^2}{2M_\pi} = \frac{140^2 - 106^2}{2 \cdot 140}$$

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1.2. Kinematika

m_1



$$e_1^* = \sqrt{(p^*)^2 + m_1^2}$$

$$p_{1,\parallel}^* = p^* \cos(\theta)$$

$$p_{1,\perp}^* = p^* \sin(\theta)$$

$$e_1^{lab} = \gamma(e_1^* + \beta p^* \cos(\theta)) \quad \gamma = \frac{E}{M} = \frac{\sqrt{P^2 + M^2}}{M} \quad \beta = \frac{P}{E} = \frac{P}{\sqrt{P^2 + M^2}}$$

$$p_{1,\parallel}^{lab} = \gamma(p^* \cos(\theta) + \beta e_1^*)$$

$$p_{1,\perp}^{lab} = p_{1,\perp}^* = p^* \sin(\theta)$$

1.3. Atom vodíku - opakování

$$\left(-\frac{(\hbar c)^2}{2m_e c^2} \Delta_e - \frac{(\hbar c)^2}{2m_p c^2} \Delta_p - \alpha \frac{\hbar c}{|\vec{r}_e - \vec{r}_p|} \right) \varphi(\vec{r}_e, \vec{r}_p) = E_{tot} \varphi(\vec{r}_e, \vec{r}_p)$$

$$\vec{R} = \frac{m_e \vec{r}_e + m_p \vec{r}_p}{m_e + m_p} \quad \vec{r} = \vec{r}_e - \vec{r}_p \quad m_{e,p} c^2 \longrightarrow m_{e,p}$$

$$\left(-\frac{(\hbar c)^2}{2(m_e + m_p)} \Delta_R - \frac{(\hbar c)^2}{2m_{ep}} \Delta_r - \alpha \frac{\hbar c}{|\vec{r}|} \right) \varphi(\vec{R}, \vec{r}) = E_{tot} \varphi(\vec{R}, \vec{r})$$

$$m_{ep} = \frac{m_e m_p}{m_e + m_p} \cong m_e; \quad \varphi(\vec{R}, \vec{r}) = \Psi(\vec{R}) \psi(\vec{r}); \quad E_{tot} = E_R + E$$

$$-\frac{(\hbar c)^2}{2(m_e + m_p)} \Delta_R \Psi(\vec{R}) = E_R \Psi(\vec{R}); \quad \left[-\frac{(\hbar c)^2}{2} \Delta_r - \alpha \frac{\hbar c}{|\vec{r}|} \right] \psi(\vec{r}) = E \psi(\vec{r})$$

1.3. Atom vodíku - opakování

$$\left(-\frac{(\hbar c)^2}{2m_{ep}} \Delta_r - \alpha \frac{\hbar c}{|\vec{r}|} \right) \psi(\vec{r}) = E \psi(\vec{r})$$

$$\psi(\vec{r}) = R(r) Y_{lm}(\theta, \varphi)$$

Laplacian: $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$
 $\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial}{\partial \vartheta} \right) + \frac{1}{r^2 \sin^2 \vartheta} \frac{\partial^2}{\partial \phi^2}$

$$\hat{L}^2 Y_{lm}(\theta, \varphi) = \hbar^2 l(l+1) Y_{lm}(\theta, \varphi); \hat{L}_z Y_{lm}(\theta, \varphi) = \hbar m Y_{lm}(\theta, \varphi)$$

$$-\frac{(\hbar c)^2}{2m_e} \frac{d^2}{dr^2} (rR(r)) - \alpha \frac{\hbar c}{r} (rR(r)) + \frac{(\hbar c)^2 l(l+1)}{2m_e r^2} (rR(r)) = E (rR(r))$$

$$u(r) = rR(r) \xrightarrow{r \rightarrow 0} r^{l+1} \quad ; \quad u(r) = rR(r) \xrightarrow{r \rightarrow \infty} 0$$

1.3. Atom vodíku - opakování

$$l=0$$

$$-\frac{(\hbar c)^2}{2m_e} \frac{d^2}{dr^2} u(r) - \alpha \frac{\hbar c}{r} u(r) = Eu(r)$$

$$u(r) = re^{-r/r_B}; \frac{d^2}{dr^2} u(r) = \frac{1}{r_B^2} u(r) - \frac{2}{rr_B} u(r)$$

$$-\frac{(\hbar c)^2}{2m_e} \frac{1}{r_B^2} u(r) + \frac{(\hbar c)^2}{2m_e} \frac{2}{rr_B} u(r) - \alpha \frac{\hbar c}{r} u(r) = Eu(r)$$

$$\frac{(\hbar c)^2}{2m_e} \frac{2}{rr_B} u(r) - \alpha \frac{\hbar c}{r} u(r) = 0; \quad -\frac{(\hbar c)^2}{2m_e} \frac{1}{r_B^2} u(r) = Eu(r)$$

$$r_B = \frac{(\hbar c)^2}{2m_e} \frac{2}{\alpha \hbar c} = \frac{\hbar c}{m_e \alpha} = \frac{197 \text{ MeV fm}}{0,511 \text{ MeV}} = 52816 \text{ fm}$$

$$E = -\frac{(\hbar c)^2}{2m_e} \frac{1}{r_B^2} = -\frac{(\hbar c)^2}{2m_e} \frac{(m_e \alpha)^2}{2} = -\frac{1}{2} m_e \alpha^2 = -\frac{1}{2} 0,511 \text{ MeV} \frac{1}{137^2} = -13,6 \text{ eV}$$

1.3. Atom vodíku - opakování

$$E = -\frac{1}{2}m_e\alpha^2 ; E = T + V = \frac{1}{2}m_e\alpha^2 - 2\frac{1}{2}m_e\alpha^2 = 13,6eV - 27,2eV$$

$$T = \frac{1}{2}m_e\left(\frac{v}{c}\right)^2 = \frac{1}{2}m_e\beta^2 ; \text{ rychlost elektronu } \beta = \frac{1}{137} = \alpha$$

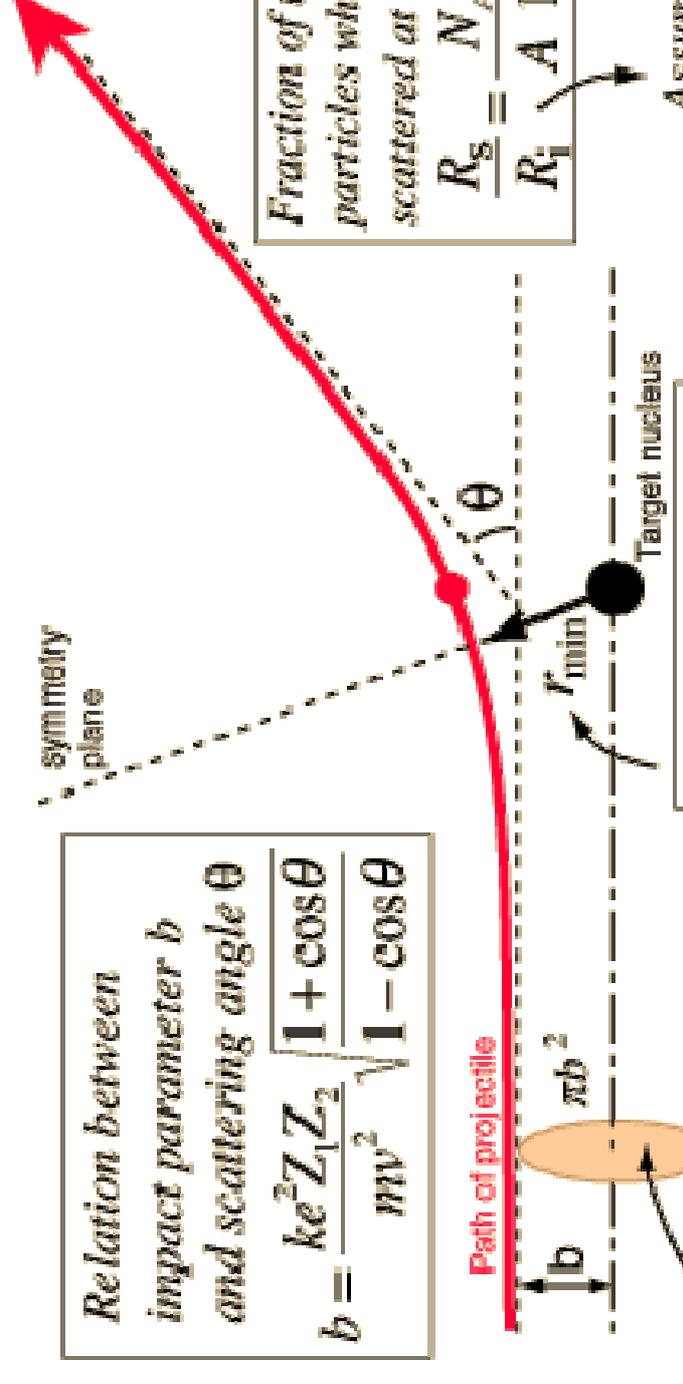
atom s Z protony :

$$E = -\frac{1}{2}m_e(Z\alpha)^2 = Z^2 13,6eV ; E_{Fe} = -26^2 13,6eV = -9,19keV$$

1.4. Rutherfordův rozptyl

Relation between impact parameter b and scattering angle θ

$$b = \frac{ke^2 Z_1 Z_2}{mv^2} \sqrt{\frac{1 + \cos\theta}{1 - \cos\theta}}$$



Relation between impact parameter b and closest approach

$$r_{\min} = \frac{b \cos\left(\frac{\theta}{2}\right)}{1 - \sin\left(\frac{\theta}{2}\right)}$$

Relation between impact parameter b and scattering angle θ

$$b = \frac{ke^2 Z_1 Z_2}{2 KE} \left(\frac{1 + \cos\theta}{1 - \cos\theta} \right)$$

Cross section for scattering above angle θ

$$\sigma = \pi b^2$$

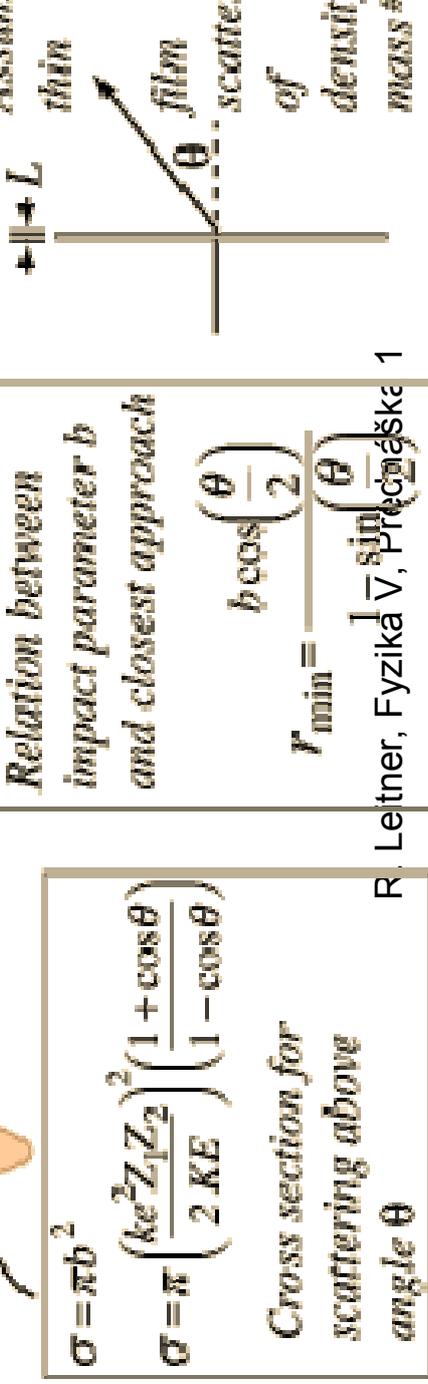
Relation between impact parameter b and scattering angle θ

$$b = \frac{ke^2 Z_1 Z_2}{mv^2} \sqrt{\frac{1 + \cos\theta}{1 - \cos\theta}}$$

Fraction of incoming particles which are scattered at angle $\geq \theta$

$$\frac{R_s}{R_i} = \frac{N_A L \rho \sigma}{A 10^{-3} \text{ kg}}$$

Assumes thin film scatterer of density ρ mass # A



1.4. Rutherfordův rozptyl klasicky - opakování

$$b(\theta) = \frac{b(\pi/2)}{\tan(\theta/2)}; b(0) \rightarrow \infty; b(\pi) \rightarrow 0$$

$$b(\pi/2) = r_{\min} / 2; \Gamma_k = \alpha Z_1 Z_2 \frac{\hbar c}{r_{\min}}; r_{\min} = \alpha Z_1 Z_2 \frac{\hbar c}{T_k}$$

$$r_{\min} = \frac{1}{137} \cdot 2 \cdot 79 \cdot \frac{197 \text{ MeV fm}}{7,7 \text{ MeV}} = 29,5 \text{ fm}$$

$$\sigma(\theta \geq \theta_0) = \pi b^2(\theta)$$

$$\sigma(\theta \geq \pi/2) = \pi b^2(\pi/2) = \pi \frac{r_{\min}^2}{4} = 3,14 \frac{(29,5 \text{ fm})^2}{4} = 683 \text{ fm}^2 =$$

$$683 \cdot 10^{-26} \text{ cm}^2 = 6,83 \cdot 10^{-24} \text{ cm}^2 = 6,83 \text{ barn}$$

1.4. Rutherfordův rozptyl klasicky - opakování

$$dP(x) = (1 - P(x))\sigma[cm^2]dx[cm]\rho[g/cm^3]N_A[1/mol] / A[g/mol]$$

$$-d(1 - P(x)) = (1 - P(x)) \frac{\sigma N_A}{A} dx$$

$$1 - P(L) = e^{-\frac{\sigma N_A L}{A}}$$

$$P(L) = 1 - e^{-\frac{\sigma N_A L}{A}} \approx 1 - \left(1 - \frac{\sigma N_A L}{A}\right) = \frac{\sigma N_A L}{A}$$

$$\sigma = 6,83b; \rho = 20g/cm^3; N_A = 6,023 \cdot 10^{23}; L = 500nm; A = 197$$

$$P(\theta > 90^\circ, 500nmAu) = \frac{6,83 \cdot 10^{-24} cm^2 \cdot 20g/cm^3 \cdot 6,023 \cdot 10^{23} mol^{-1}}{197g/mol} = 500 \cdot 10^{-7} cm =$$

$$2,09 \cdot 10^{-5} \approx \frac{1}{50000}$$

1.4. Rutherfordův rozptyl klasicky - opakování

Diferenciální účinný průřez

$$d\sigma = 2\pi b(\theta)db = 2\pi b(\theta)\left|\frac{db}{d\theta}\right|d\theta$$

$$b(\theta) = \frac{b(\pi/2)}{\tan(\theta/2)}; \frac{db}{d\theta} = -\frac{b(\pi/2)}{\tan^2(\theta/2)} \frac{1}{2}$$

$$d\sigma = -2\pi \frac{b(\pi/2)}{\tan(\theta/2)} \frac{b(\pi/2)}{\tan^2(\theta/2)} \frac{1}{2} d\theta$$

$$\frac{d\sigma}{d\theta} = -\frac{\pi(r_{\min}/2)^2}{\sin^3(\theta/2)} \cos(\theta/2) = -\frac{1}{2} \frac{\pi(r_{\min}/2)^2}{\sin^4(\theta/2)} 2\sin(\theta/2)\cos(\theta/2) = -\frac{1}{2} \frac{\pi(r_{\min}/2)^2}{\sin^4(\theta/2)} \sin(\theta)$$

$$\frac{d\sigma}{-\sin(\theta)d\theta} = \frac{d\sigma}{d\cos(\theta)} = \frac{\pi(r_{\min}^2/8)}{\sin^4(\theta/2)}$$

1.4. Výsledek Rutherfordova pokusu

- Částice alpha se rozptylují jakoby jádro bylo bodové
- Rozměr jádra je tudíž menší než $R_{\min}=30 \text{ fm}$, tj. asi 2000 krát menší než rozměr atomů
- Alfa částice pronikaly do lehkých jader (Be) a tudíž jádro má nejspíš konečný rozměr.

1.5. Definice účinného průřezu v kvantové mechanice

$$d\sigma = \frac{\vec{J}_{out} \cdot d\vec{S}}{J_{in}}; \vec{j} = \frac{\hbar}{i2m} (\psi^* \vec{\nabla} \psi - \psi \vec{\nabla} \psi^*);$$

$$\psi_{in} = e^{\frac{ipz}{\hbar}}; j_{in,z} = \frac{\hbar}{i2m} (e^{-i\frac{pz}{\hbar}} i \frac{p}{\hbar} e^{\frac{ipz}{\hbar}} - e^{\frac{ipz}{\hbar}} -i \frac{p}{\hbar} e^{-i\frac{pz}{\hbar}}) = \frac{\hbar}{i2m} 2i \frac{p}{\hbar} = \frac{p}{mc^2} = \frac{p}{m}$$

$$\psi_{out} = f(\theta) \frac{e^{\frac{ipr}{\hbar}}}{r}; j_{out,x} = \frac{\hbar}{i2m} (f^*(\theta) \frac{e^{-i\frac{pr}{\hbar}}}{r} - i \frac{p}{\hbar} f(\theta) \frac{e^{\frac{ipr}{\hbar}}}{r} \frac{\partial r}{\partial x} - f(\theta) \frac{e^{\frac{ipr}{\hbar}}}{r} (-i \frac{p}{\hbar}) f^*(\theta) \frac{e^{-i\frac{pr}{\hbar}}}{r} \frac{\partial r}{\partial x}) =$$

$$\frac{\hbar}{i2m} f^*(\theta) f(\theta) 2i \frac{p}{\hbar} \frac{x}{r} = \frac{p}{mc^2} |f(\theta)|^2 \frac{x}{r}$$

$$\vec{j}_{out} = \frac{p}{m} \frac{|f(\theta)|^2}{r^2} \vec{r}$$

$$d\sigma = \frac{\vec{J}_{out} \cdot d\vec{S}}{J_{in}} = \frac{p}{m} \frac{|f(\theta)|^2}{r^2} \frac{\vec{r}}{r} \cdot \vec{r}^2 d\Omega \frac{r}{r} = |f(\theta)|^2 d\Omega$$

$$\frac{d\sigma}{d \cos(\theta) d\phi} = |f(\theta)|^2$$

1.5. Rozptyl v kvantové mechanice

$$\begin{aligned}
 \left(\Delta + \left(\frac{\sqrt{2ME}}{\hbar c} \right)^2 \right) \psi^{(0)}(\vec{r}) &= 0 \quad ; \quad \psi^{(0)}(\vec{r}) = e^{+ikz} \quad ; \quad k = \frac{\sqrt{2ME}}{\hbar c} \\
 \left(-\frac{(\hbar c)^2}{2M} \Delta + V(\vec{r}) \right) \psi(\vec{r}) &= E \cdot \psi(\vec{r}) \quad ; \quad \Delta + \left(\frac{\sqrt{2ME}}{\hbar c} \right)^2 \psi(\vec{r}) = \frac{2M}{(\hbar c)^2} V(\vec{r}) \cdot \psi(\vec{r}) = \frac{2M}{(\hbar c)^2} V(\vec{r}) \cdot \psi(\vec{r}) \\
 (\Delta + k^2) \psi^{(n)}(\vec{r}) &= \frac{2M}{(\hbar c)^2} V(\vec{r}) \cdot \psi^{(n-1)}(\vec{r}) \quad ; \quad \psi^{(n)}(\vec{r}) = -\frac{2M}{(\hbar c)^2} \int V(\vec{r}') \cdot \psi^{(n-1)}(\vec{r}') \cdot G(\vec{r} - \vec{r}') \cdot d^3 \vec{r}' \\
 (\Delta + k^2) G(\vec{r} - \vec{r}') &= \delta(\vec{r} - \vec{r}') \quad ; \quad G(\vec{r} - \vec{r}') = \frac{1}{4\pi} \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \\
 \psi^{(1)}(\vec{r}) &= -\frac{M}{2\pi(\hbar c)^2} \int V(\vec{r}') \cdot \psi^{(0)}(\vec{r}') \cdot \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \cdot d^3 \vec{r}' + \psi^{(0)}(\vec{r}) \\
 \psi^{(2)}(\vec{r}) &= -\frac{M}{2\pi(\hbar c)^2} \int V(\vec{r}') \cdot \psi^{(1)}(\vec{r}') \cdot \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \cdot d^3 \vec{r}' = -\frac{M}{2\pi(\hbar c)^2} \int V(\vec{r}') \cdot \left(-\frac{M}{2\pi(\hbar c)^2} \int V(\vec{r}'') \cdot \frac{e^{ik|\vec{r} - \vec{r}''|}}{|\vec{r} - \vec{r}''|} \cdot d^3 \vec{r}'' + \psi^{(0)}(\vec{r}') \right) \cdot \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \cdot d^3 \vec{r}' + \psi^{(0)}(\vec{r}) = \\
 &= -\frac{M}{2\pi(\hbar c)^2} \int V(\vec{r}') \cdot \left(-\frac{M}{2\pi(\hbar c)^2} \int V(\vec{r}'') \cdot \frac{1}{4\pi} \cdot \frac{e^{ik|\vec{r} - \vec{r}''|}}{|\vec{r} - \vec{r}''|} \cdot d^3 \vec{r}'' \right) \cdot \frac{1}{4\pi} \cdot \frac{e^{ik|\vec{r} - \vec{r}'|}}{|\vec{r} - \vec{r}'|} \cdot d^3 \vec{r}' + \psi^{(0)}(\vec{r})
 \end{aligned}$$

1.5. Rozptyl v kvantové mechanice

$$|\mathbf{r} - \mathbf{r}'| = \sqrt{(\vec{r} - \vec{r}') \cdot (\vec{r} - \vec{r}')} \cong \sqrt{rr - 2rr' \cos\theta} = \sqrt{rr(1 - 2r'r'/rr)}$$

$$\cong \sqrt{rr(1 - r'r'/rr)} = r - r'n$$

$$\frac{e^{ik|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} \cong \frac{e^{ikr} e^{-ikr'n}}{r(1-r'n/r)} \cong \frac{e^{ikr}}{r} \cdot e^{-ikr'n}$$

$$\psi^{(1)}(\mathbf{r}) = \left(-\frac{M}{2\pi(\hbar c)^2} \int e^{-ik\vec{r}'\cdot\vec{n}} \cdot V(\mathbf{r}') \cdot e^{ikz'} \cdot d^3r' \right) \frac{e^{+ikr}}{r} = f(\theta) \frac{e^{+ikr}}{r}$$

$$f(\theta, \phi) = f^{(1)}(\theta, \phi) = -\frac{M}{2\pi(\hbar c)^2} \int e^{-ik\vec{r}'\cdot\vec{n}} \cdot V(\mathbf{r}') \cdot e^{ikz'} \cdot d^3r'$$

1.5. Rozptyl v kvantové mechanice

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\cos(\theta)d\phi} = \left| -\frac{M}{2\pi(\hbar c)^2} \int e^{-i(\vec{k}_f - \vec{k}_i)\vec{r}} \cdot V(r') \cdot d^3r' \right|^2 =$$

$$\frac{M^2}{4\pi^2(\hbar c)^4} \left| \int e^{-i\frac{(\vec{p}_f - \vec{p}_i)\vec{r}}{\hbar c}} \cdot V(r') \cdot d^3r' \right|^2 = \frac{M^2}{4\pi^2(\hbar c)^4} \left| \int e^{-i\frac{\vec{q}\vec{r}}{\hbar c}} \cdot V(r') \cdot d^3r' \right|^2$$

$$\left[\frac{d\sigma}{d\Omega} \right] = \frac{\text{GeV}^2}{\text{GeV}^4 \text{fm}^4} (\text{GeV} \cdot \text{fm}^3)^2 = \text{fm}^2$$

$$|f(\theta)|^2 = \frac{M^2}{4\pi^2(\hbar c)^4} \left| \int e^{-i\frac{\vec{q}\vec{r}}{\hbar c}} \cdot V(r') \cdot d^3r' \right|^2$$